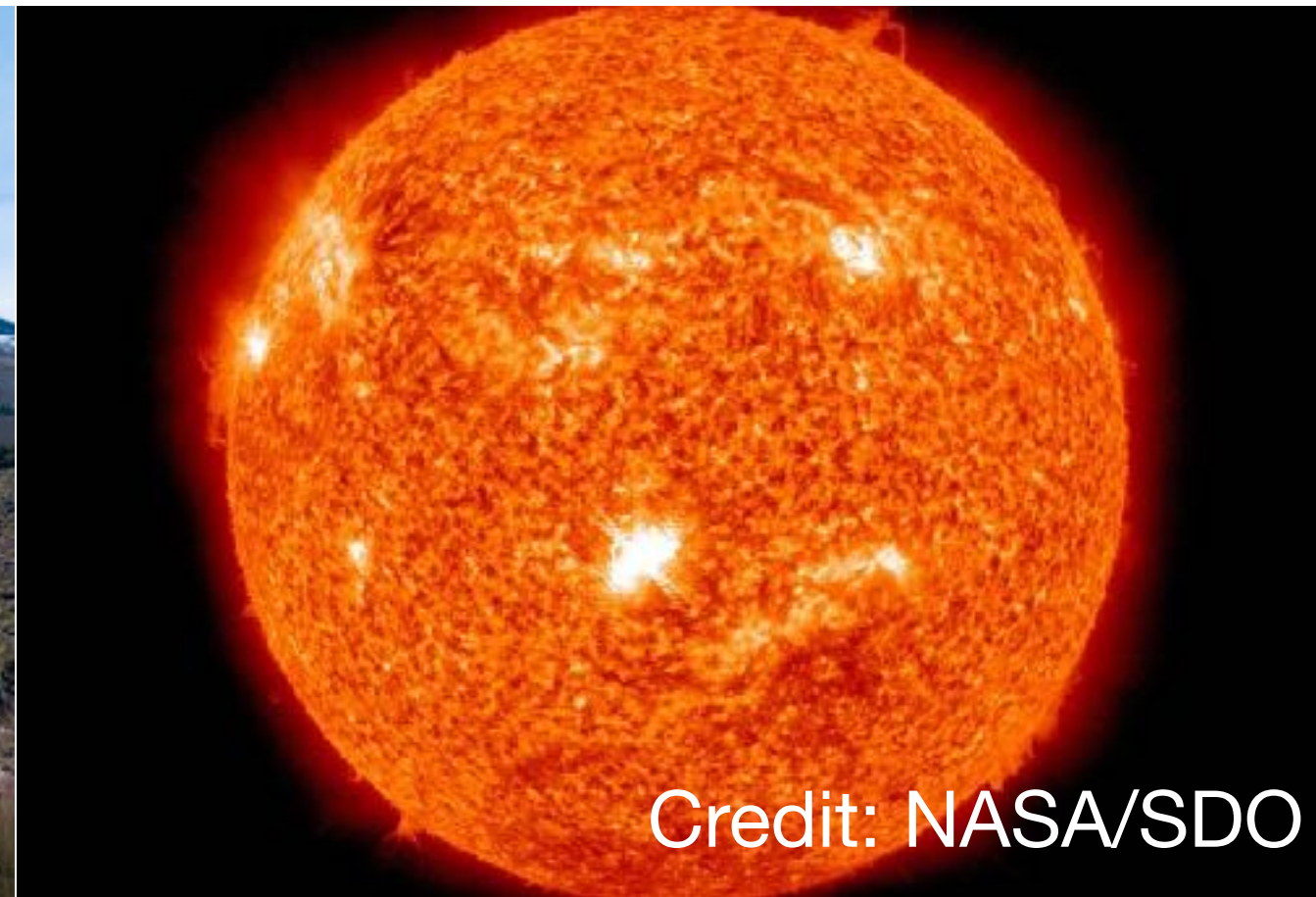


Small-Scale Magnetic Fields are Critical to Shaping Continual Solar Gamma-Ray Emission



HAWC



Credit: NASA/SDO



Fermi-LAT

Jung-Tsung Li
(Ohio State University)

CCAPP Symposium 2023

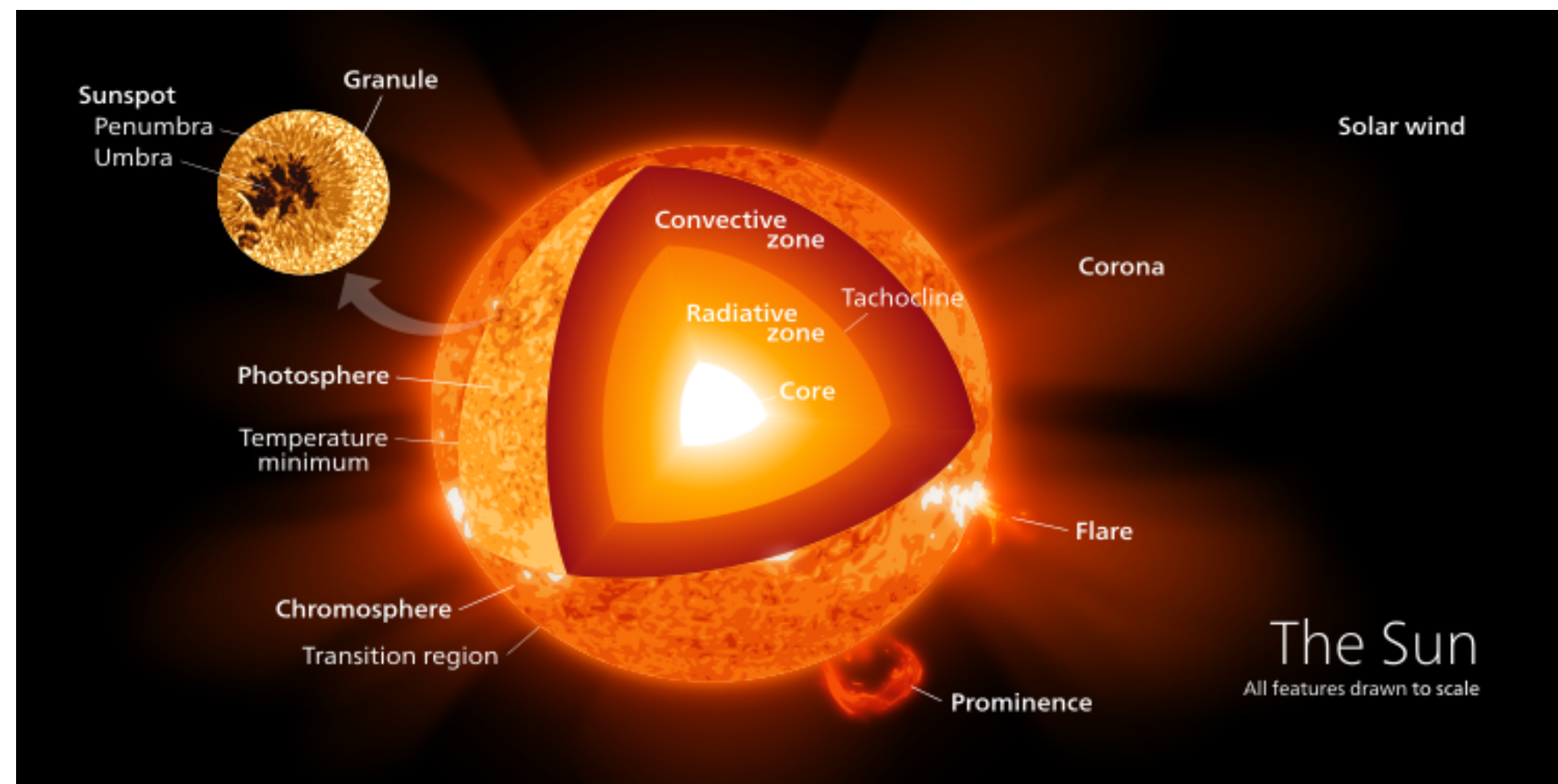


Collaborators: John Beacom, Spencer Griffith, and Annika Peter



Why is *continual* solar gamma-ray emission interesting?

(Because the Sun itself doesn't emit continual gamma rays)



- Photosphere temperature is **6000 Kelvin** — visible light (~ 1 eV)
- Corona temperature can reach as high as **4 million Kelvin**
 - **EUV and X-ray** ($\lesssim 1$ keV)
 - Heating due to wave-driven turbulence and reconnection
- Solar flare and coronal mass ejection emit gamma rays up to **few GeV**
 - Due to non-thermal particle acceleration from shock-like structures
 - Signals are transient — can be removed out from continual emission

Continual gamma rays from solar halo

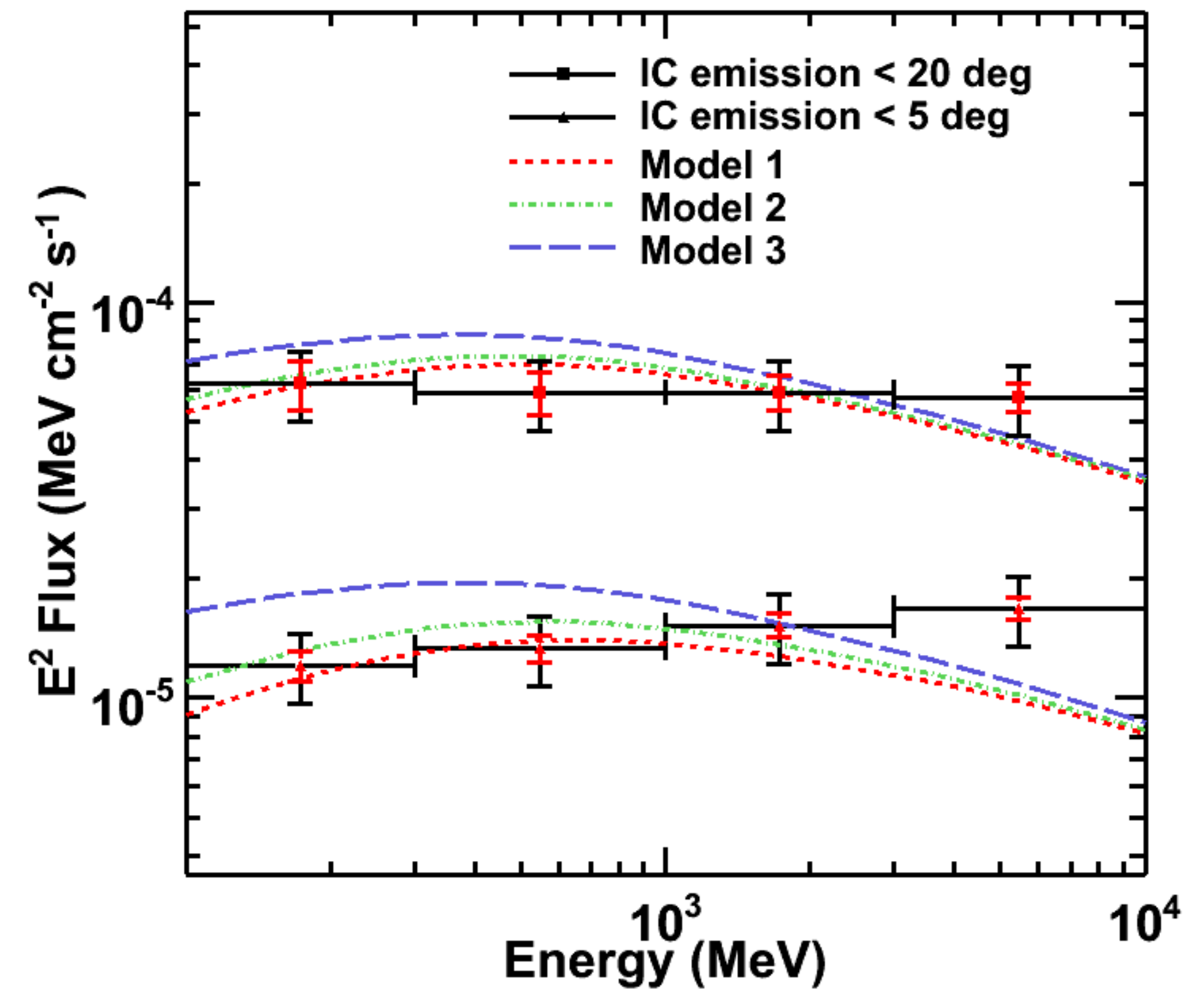
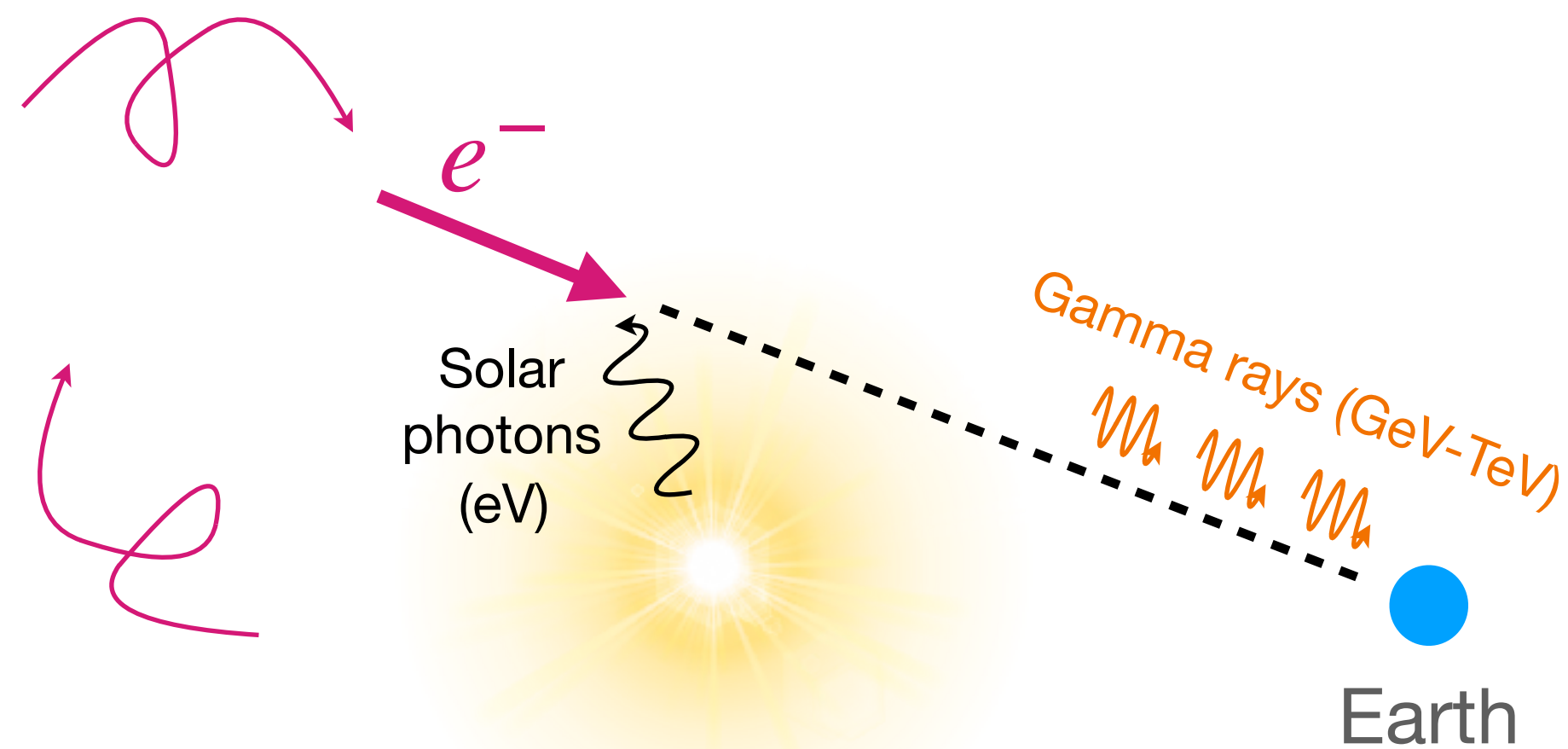
(Not the focus of this talk)

Inverse-Compton scattering in the solar halo

$$e^- + \gamma \rightarrow e^- + \gamma$$

See Moskalenko, Porter & Diego 2006;
Orlando & Strong 2007;
Abdo et al 2011

Galactic cosmic-ray electron

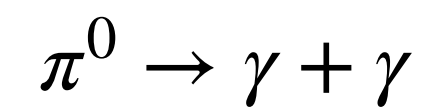
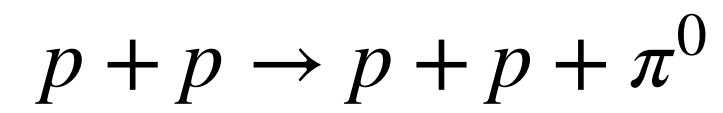


(Fermi Collaboration; Abdo et al 2011)

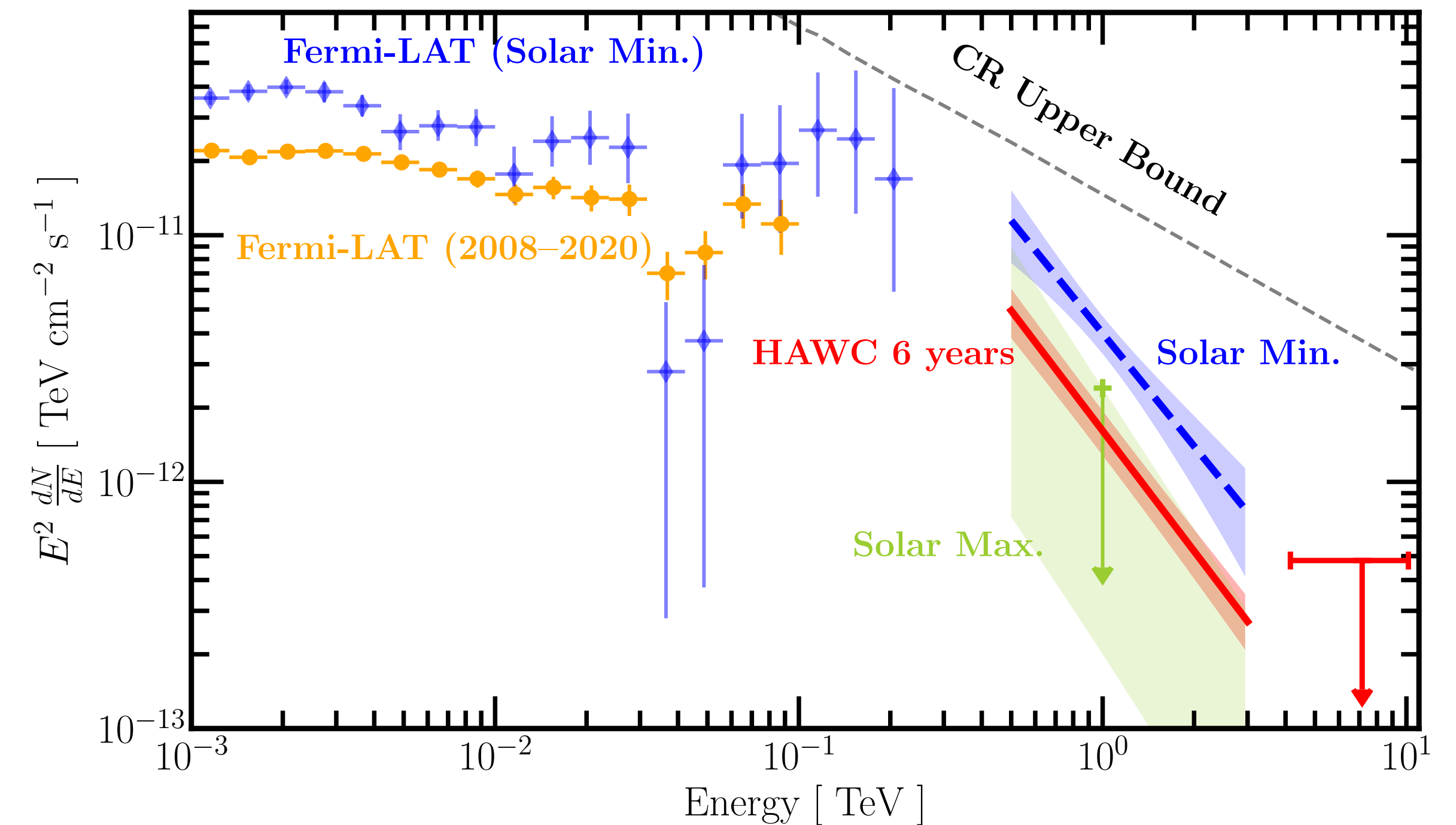
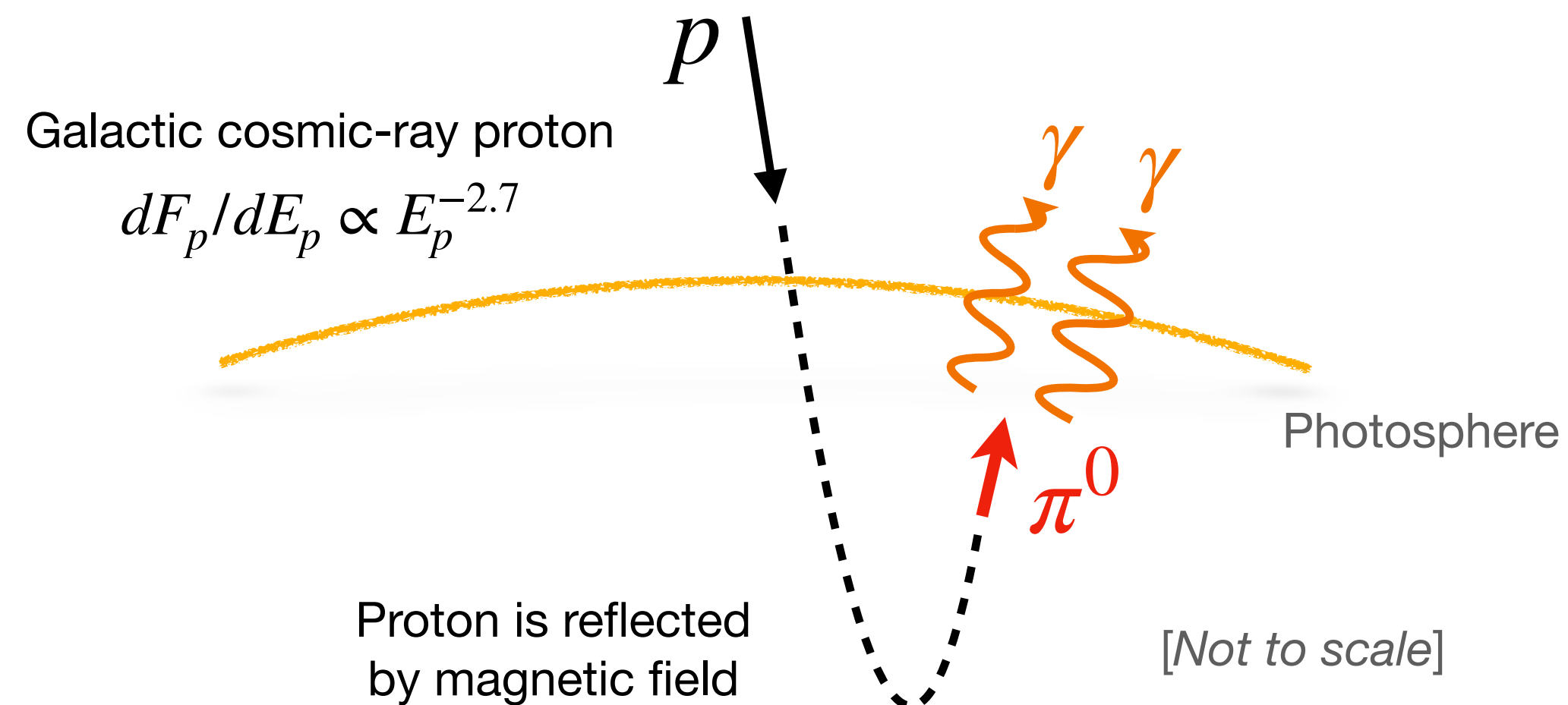
Continual gamma rays from solar disk

Focus of this talk!

Hadronic scattering in the solar disk

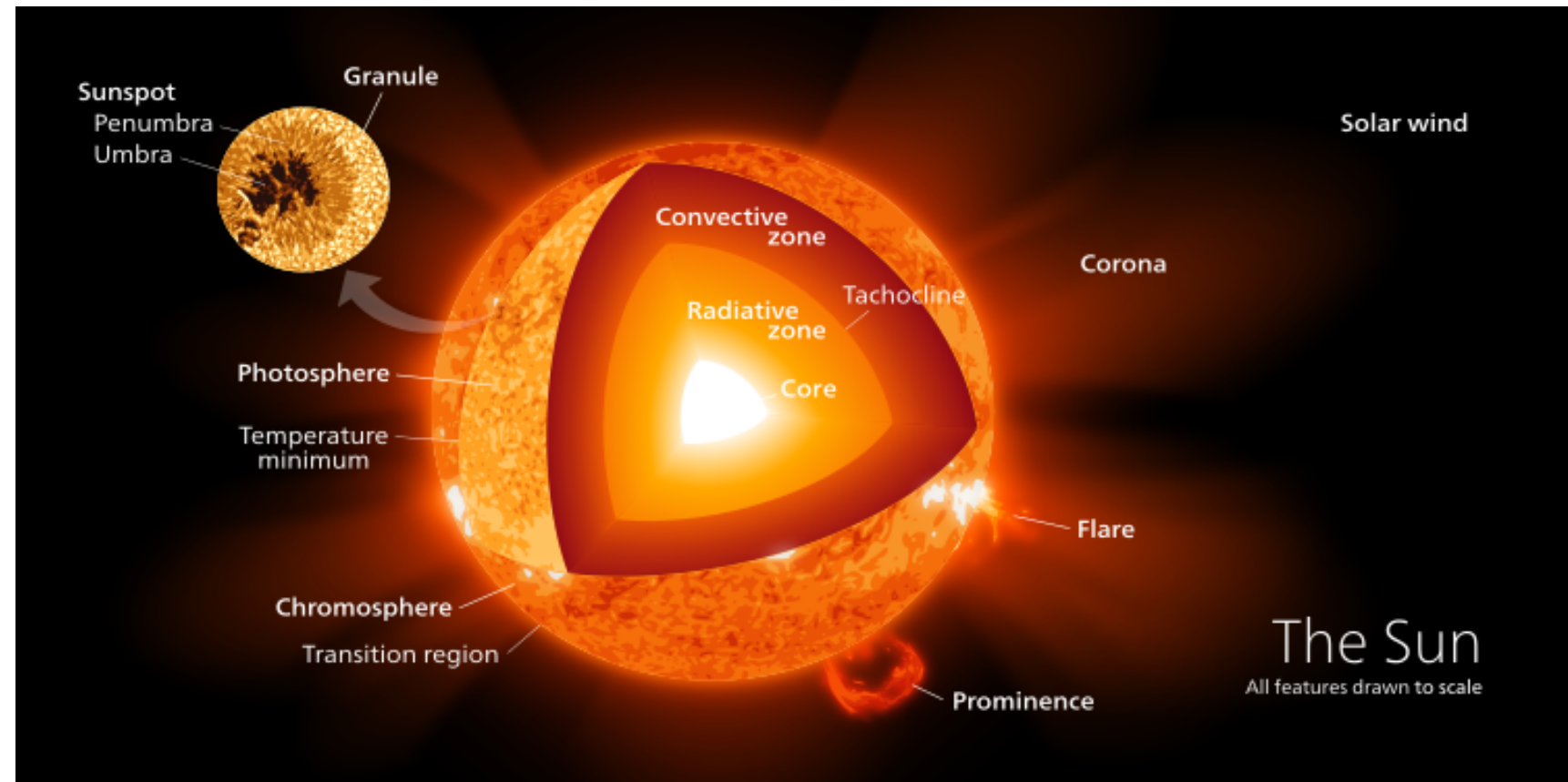


(See Seckel, Stanev & Gaisser 1991)



(HAWC Collaboration; Albert et al 2023)

Theoretical challenges for solar disk emission



1. Magnetic field structures determining the observed gamma-ray spectrum

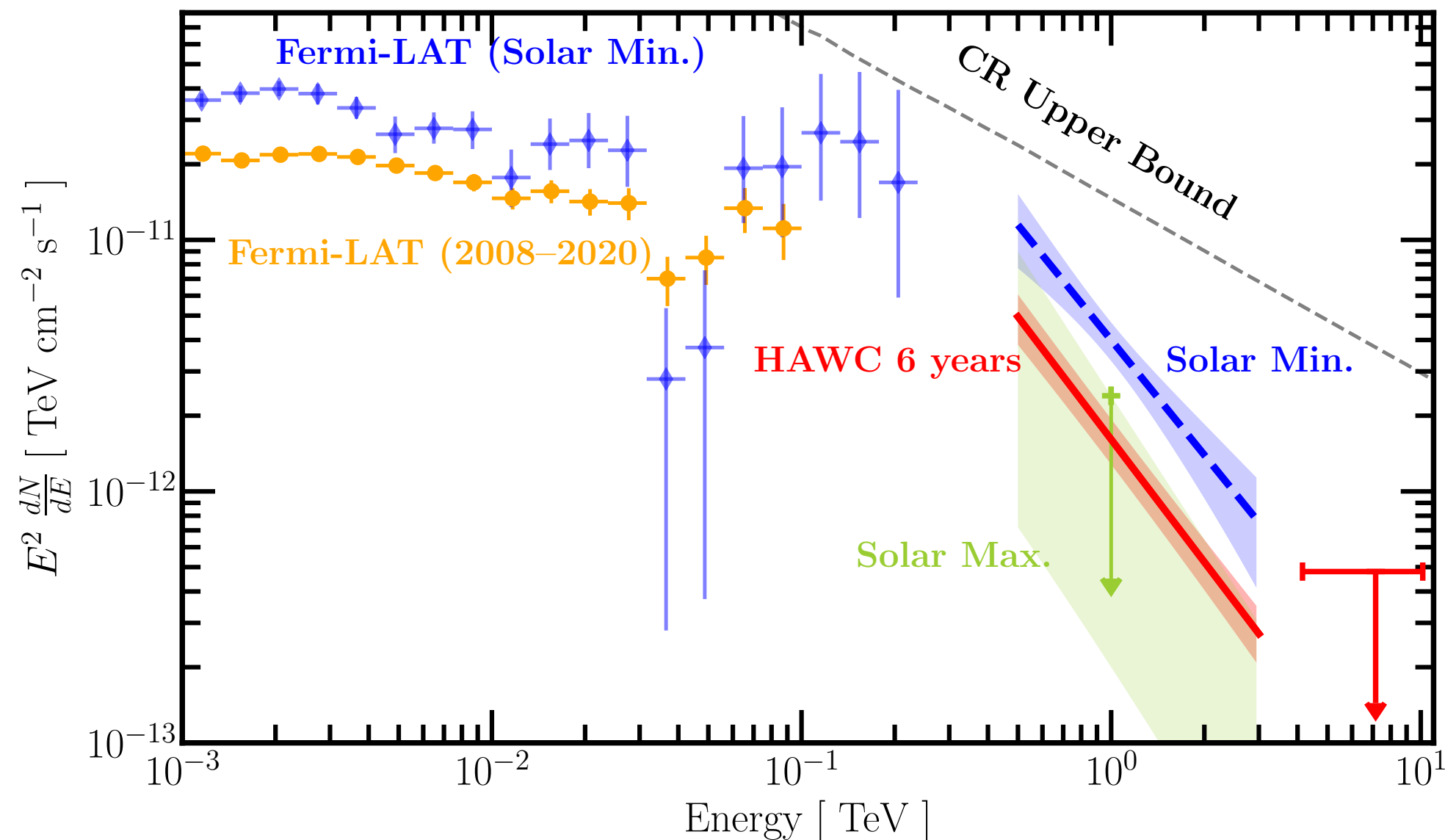
- Solar magnetic field is **multi-scale**. How do we think this problem?

2. Spectral shape

- Hard spectrum for $\lesssim 200$ GeV ($dN_\gamma/dE_\gamma \sim E_\gamma^{-2.2}$)
- Soft spectrum at ~ 1 TeV ($dN_\gamma/dE_\gamma \sim E_\gamma^{-3.6}$)

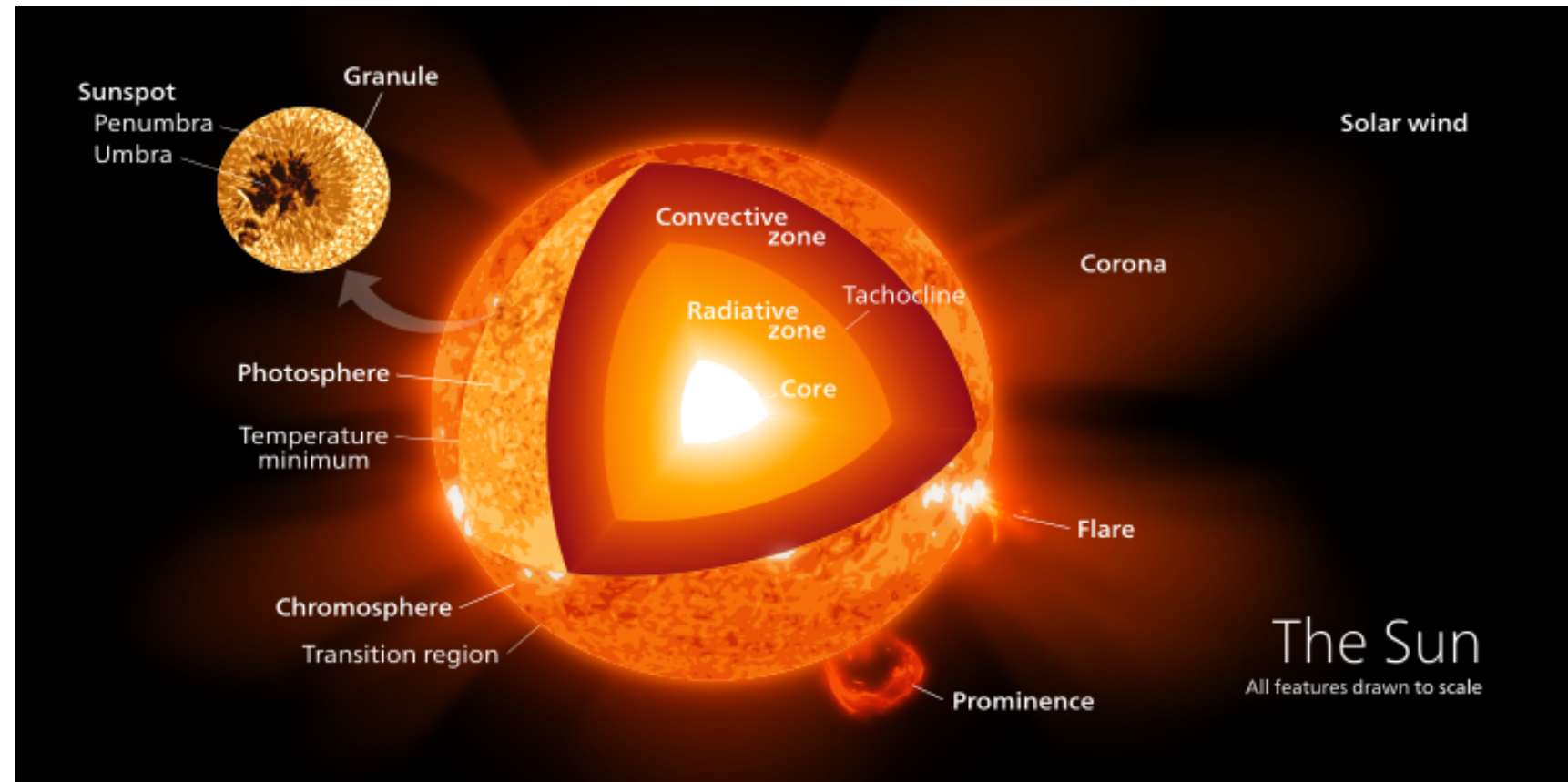
3. Gamma-ray emission anti-correlated with solar activity

- Higher gamma-ray flux at solar min
- GCR Transport? Active region activity? Small-scale convection at quiet photosphere?



(HAWC Collaboration; Albert et al 2023)

Theoretical challenges for solar disk emission



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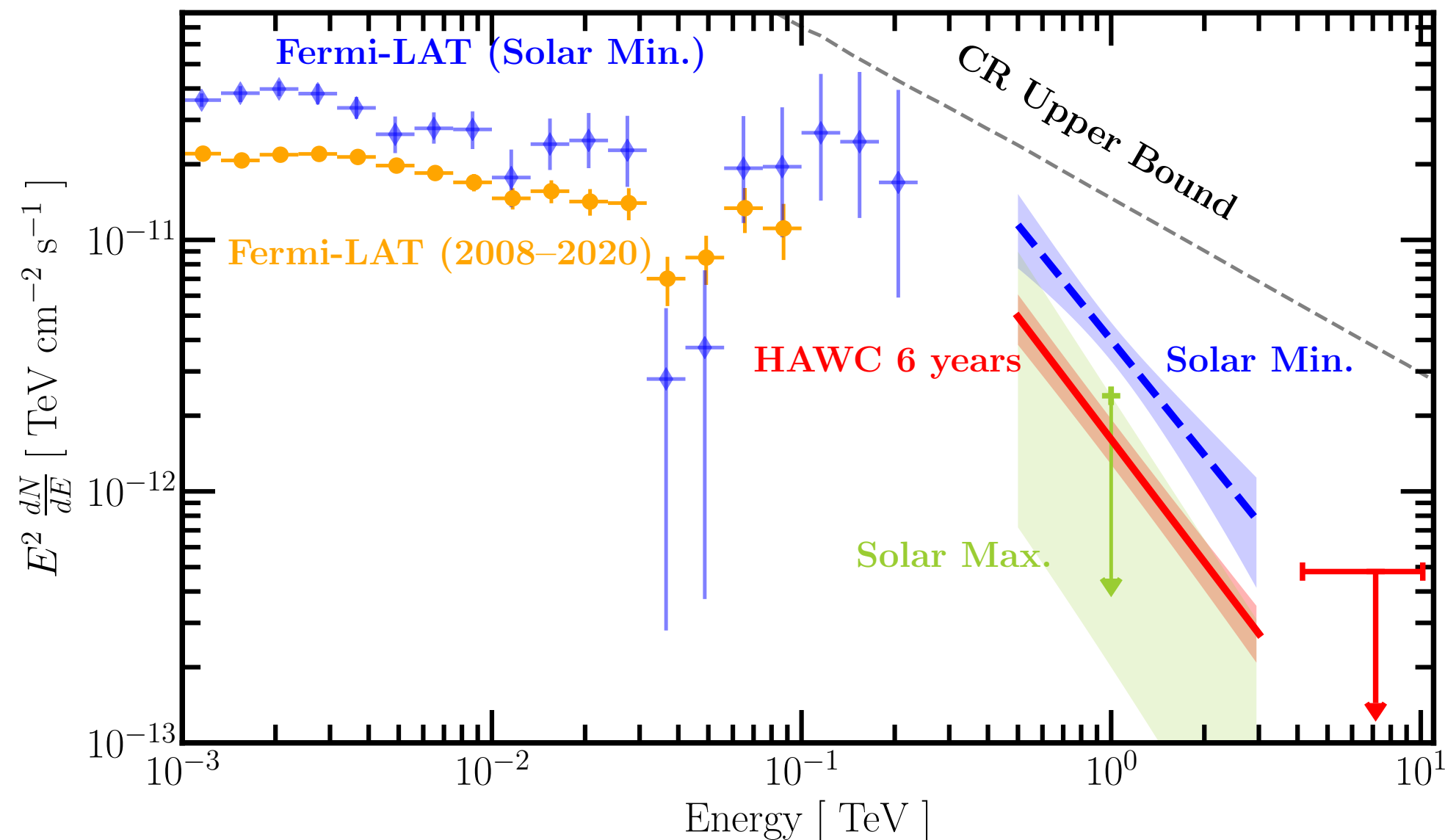
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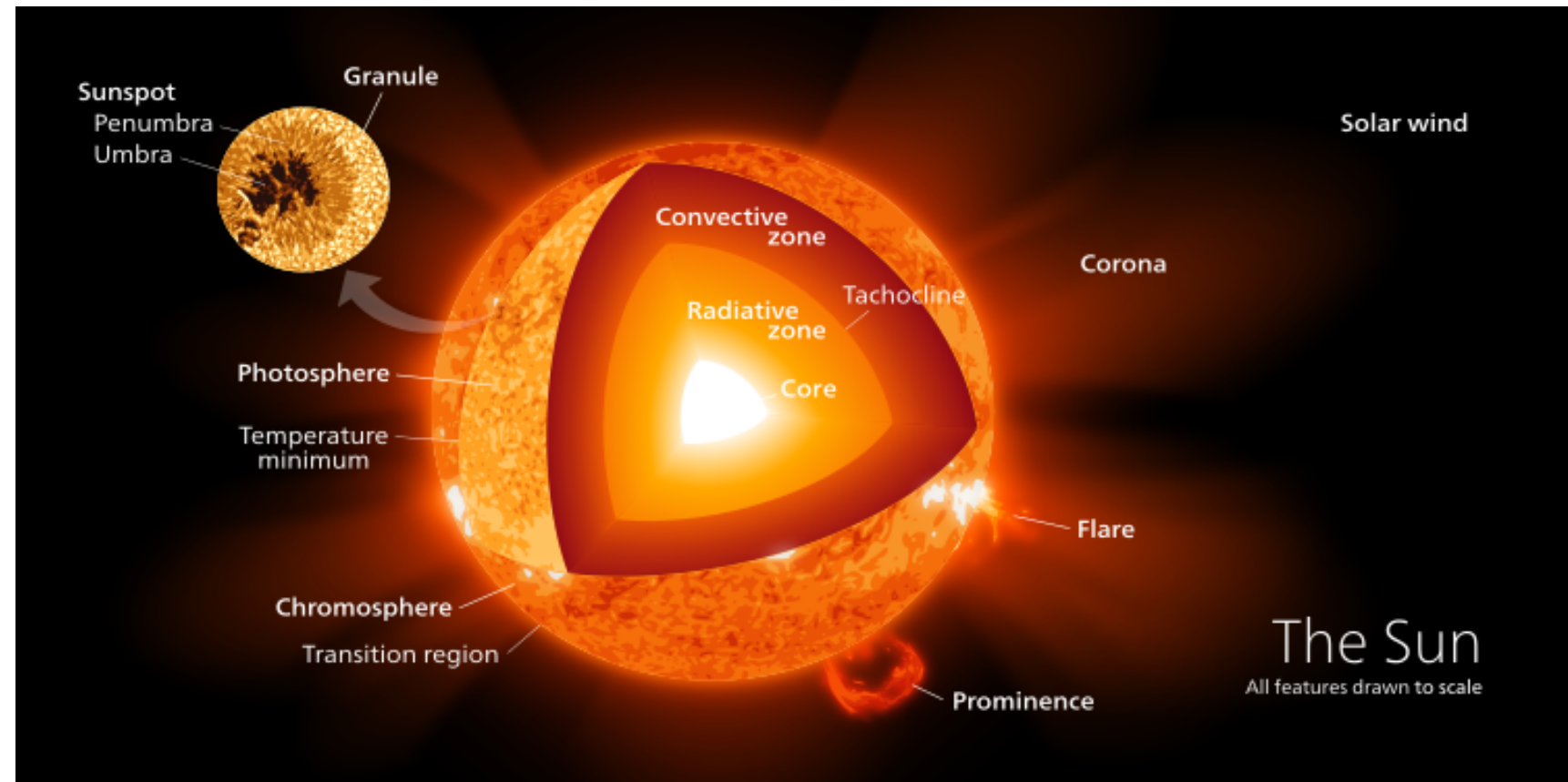
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Theoretical challenges for solar disk emission



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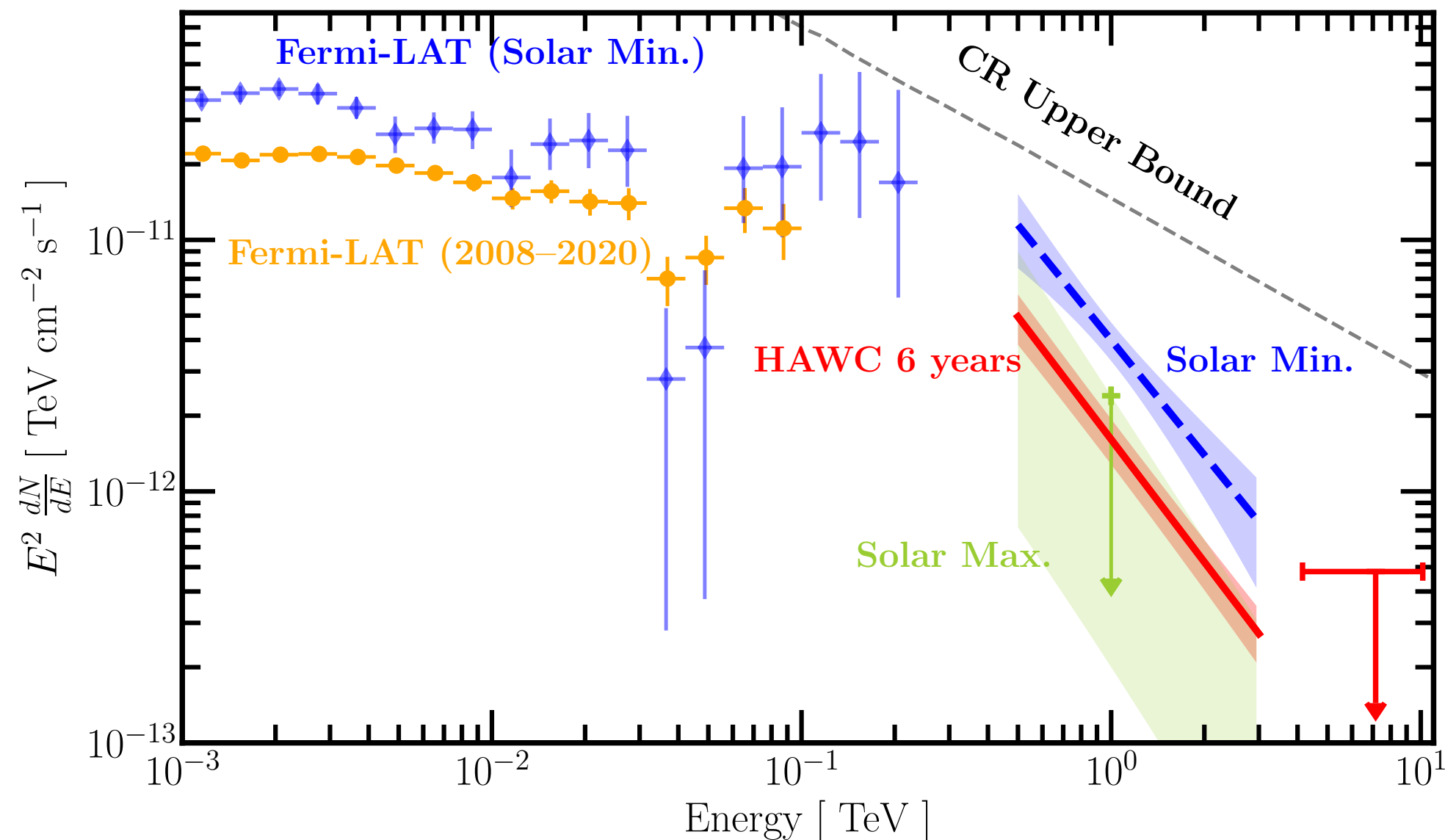
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(HAWC Collaboration; Albert et al 2023)

The Sun's magnetic structure is complex

- It is **impractical** to consider all structures **at all scales** in one study
- The goal is understand the nature of the problem: What critical magnetic structures should we consider?

In this work

- We consider quiet region of the Sun that forms the network field and open magnetic field lines
- Open field lines extends to interplanetary space and become the interplanetary magnetic fields

Stage 1: Solar Modulation

Using Force-Field Model

- Full cosmic ray transport equation, in the solar system frame (Parker 1965; Gleeson & Webb 1978)

$$\frac{\partial U_p}{\partial t} + \nabla \cdot (C \mathbf{V}_{sw} U_p) - \nabla \cdot (\kappa \cdot \nabla U_p) + v_D \cdot \nabla U_p + \frac{1}{3} \frac{\partial}{\partial p} (p \mathbf{V}_{sw} \cdot \nabla U_p) = 0$$

Rate change
Convection
Diffusion
Drift
Momentum loss

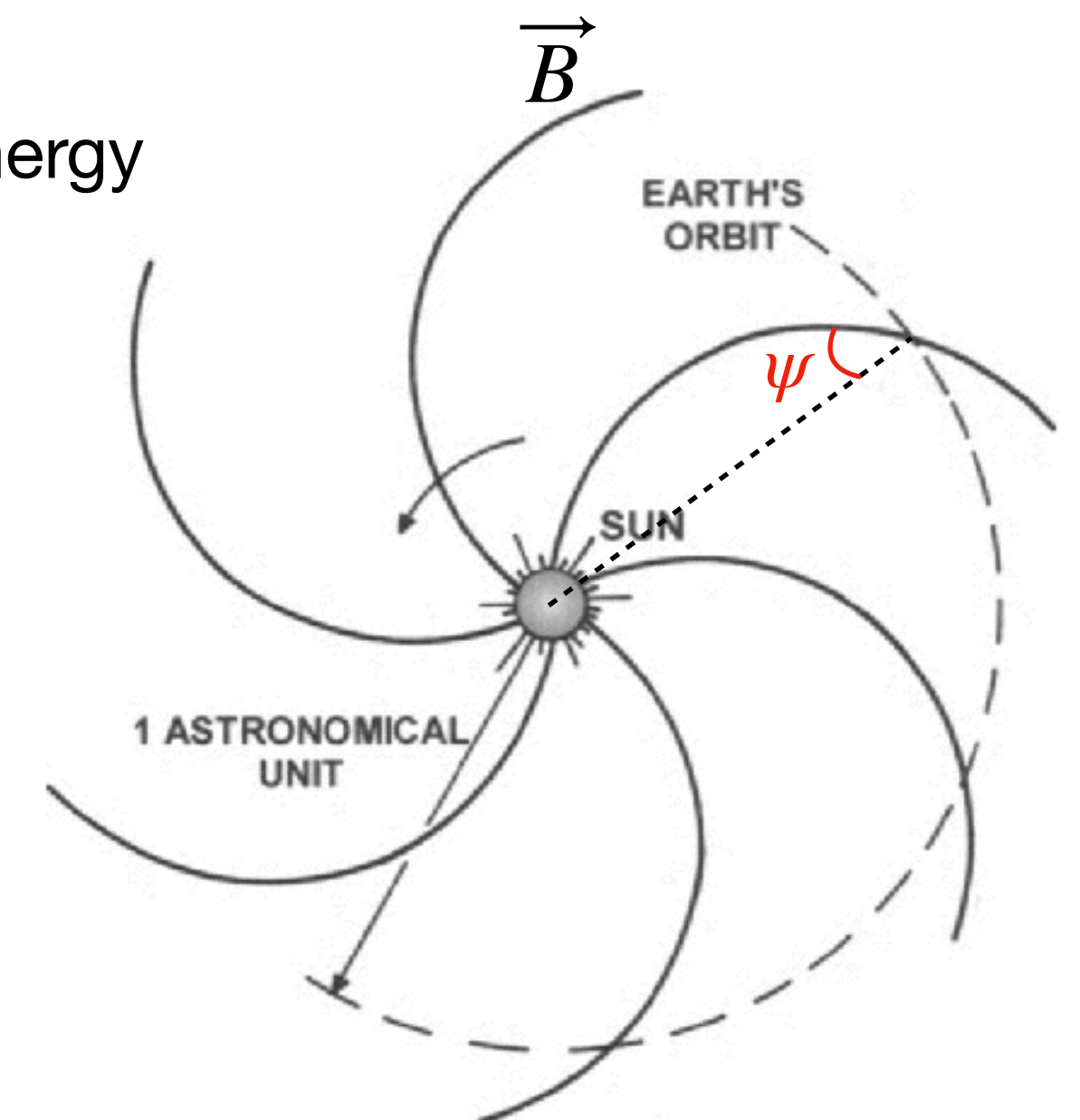
- 1D force-field model: convection flux balances diffusion flux (Gleeson & Axford 1966)

1. Force-field solution $\frac{J_E(E, r_1)}{E^2 - E_0^2} = \frac{J_E(E + \Delta\Phi, r_2)}{(E + \Delta\Phi)^2 - E_0^2}$ where $\Delta\Phi$ is modulation potential energy

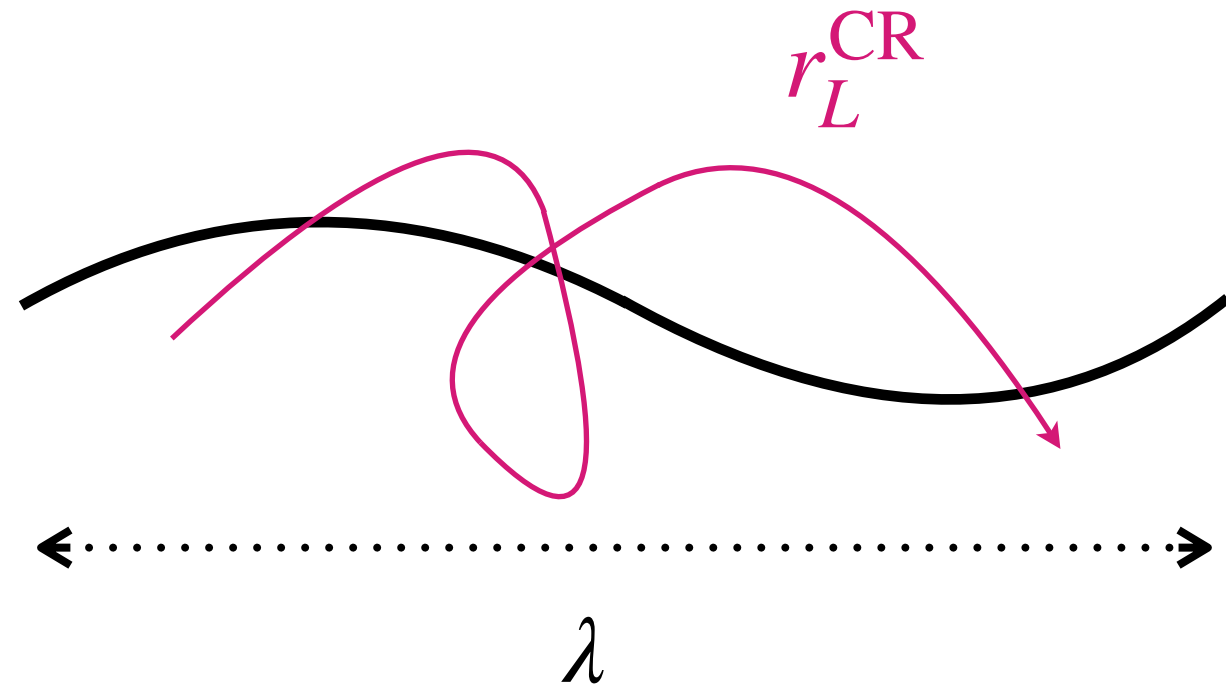
2. Characteristic eqn $\frac{dE}{dr} = \frac{V_{sw}}{3\kappa_{rr}} \frac{(E^2 - E_0^2)}{E}$

$\kappa_{rr} = \kappa_{\parallel} \cos^2 \psi + \kappa_{\perp} \sin^2 \psi$ in the plane, with $\kappa_{\parallel} \gg \kappa_{\perp}$ in the inner heliosphere

κ_{\parallel} is determined from CR resonant interaction with magnetic turbulence



Quasi-Linear Theory (QLT)



- Quasi-linear theory describes the slow evolution of the particle distribution in a weak turbulent plasma back to a marginally stable state.

$$\kappa_{\parallel} = \frac{v^2}{4} \int_{\mu_{\min,s}}^1 \frac{(1 - \mu^2)^2}{D_{\mu\mu}} d\mu \quad D_{\mu\mu} = \frac{1 - \mu^2}{2|\mu|v} \left(\frac{\Omega_{0,s}}{|\langle \mathbf{B} \rangle|} \right)^2 V_{\text{sw}}(r) E_{\text{B},xx}(f_{\text{res}}, r)$$

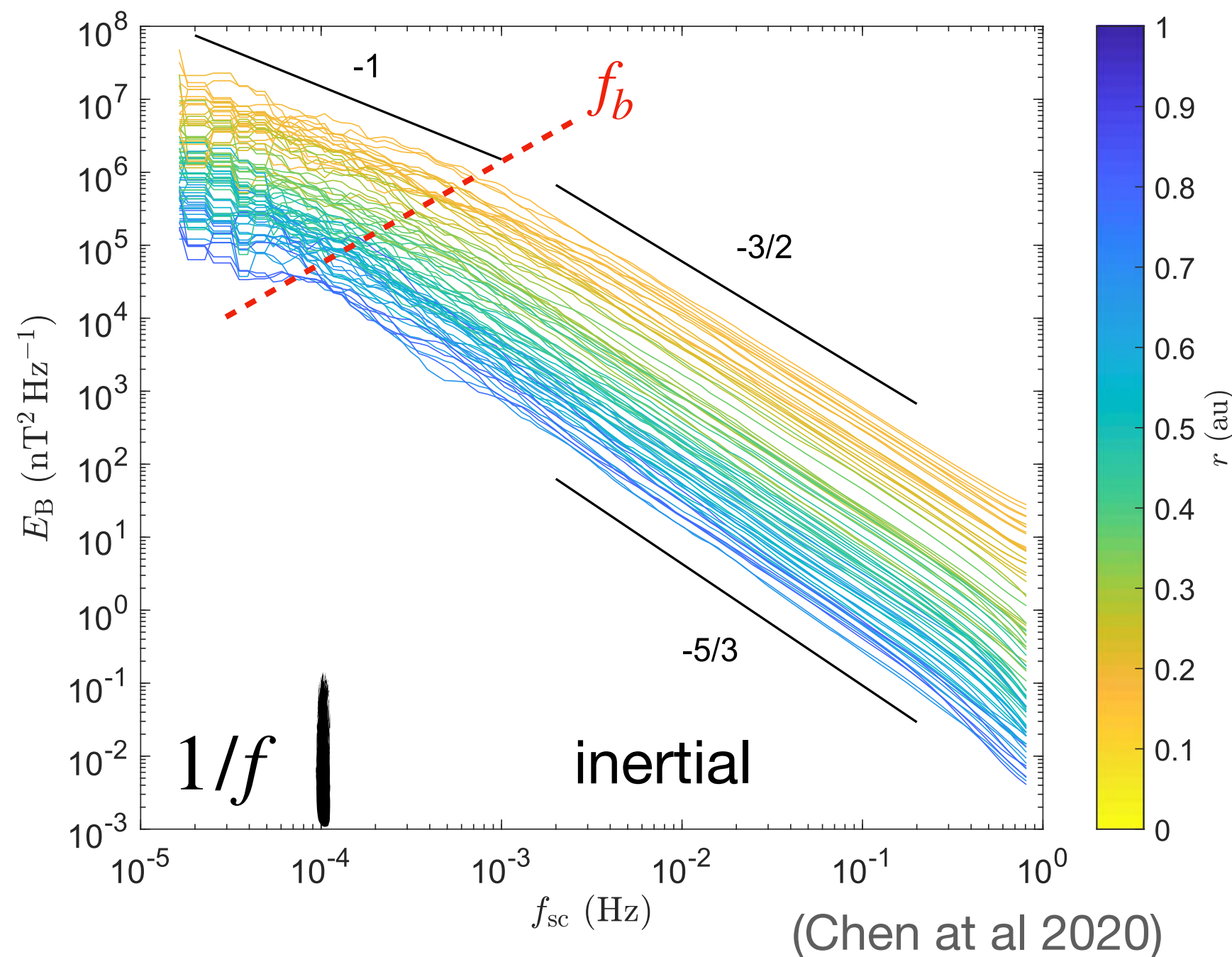
(Jokipii 1966)

μ : cosine of pitch angle

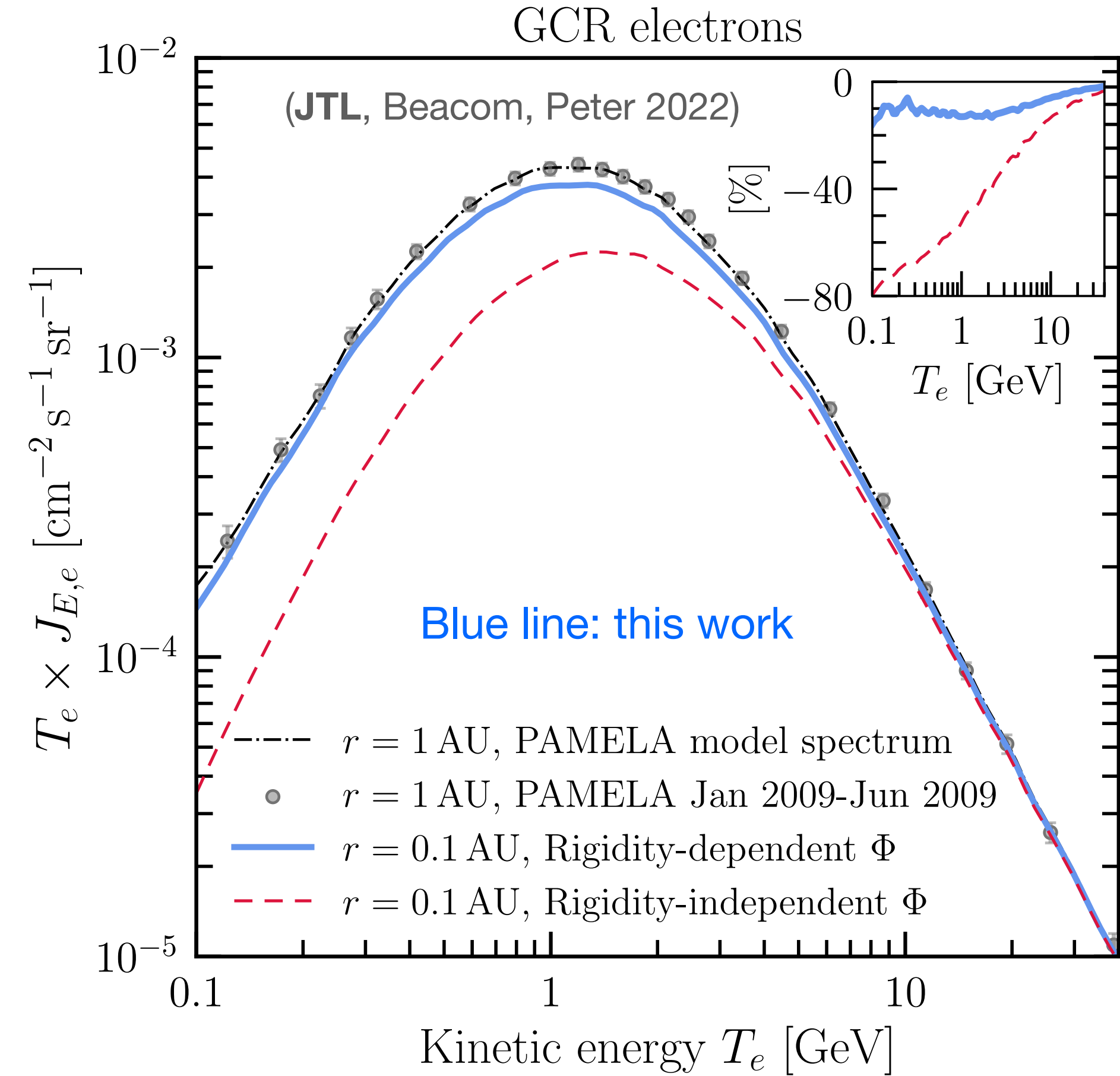
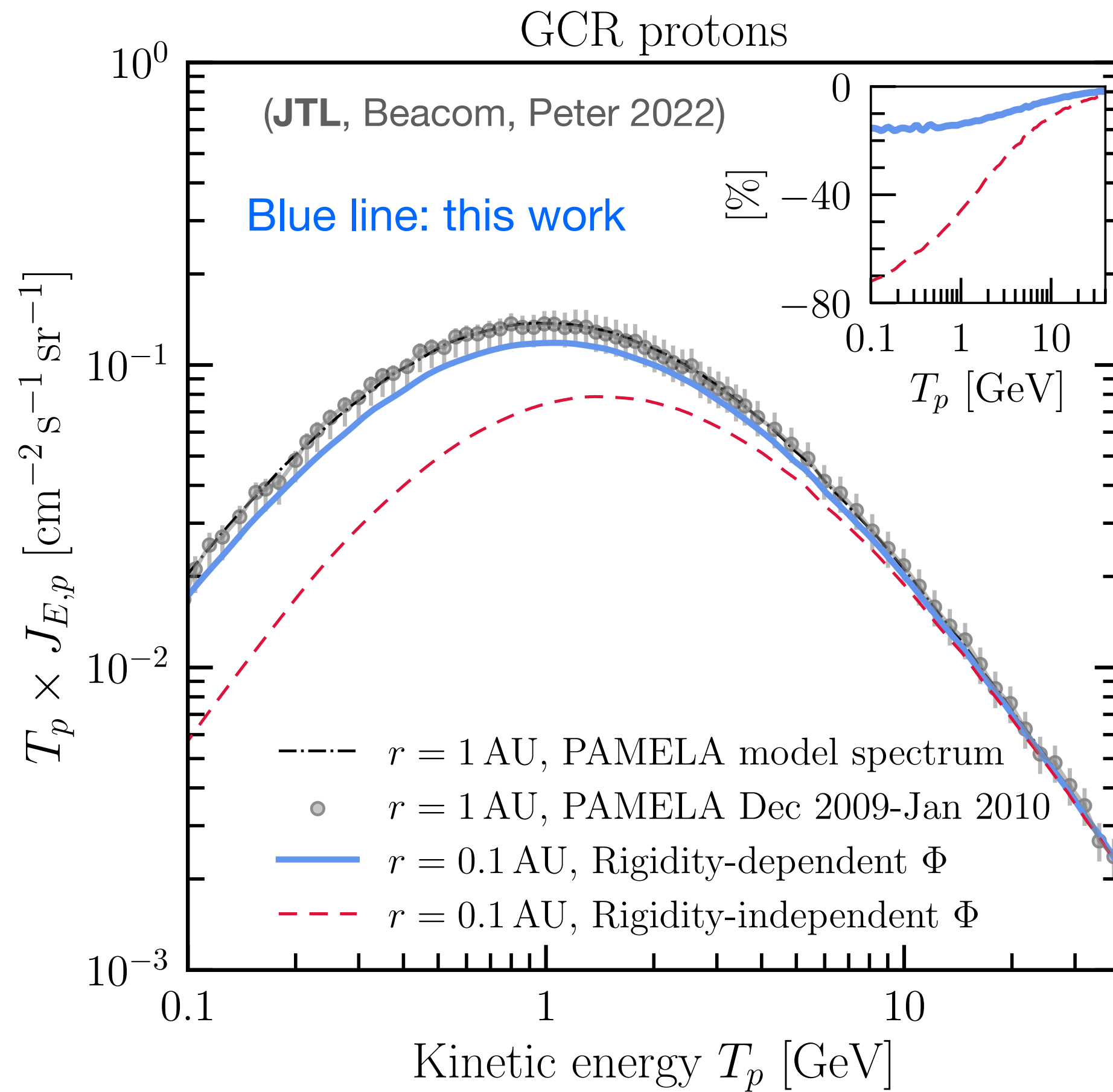
E_{B} : magnetic power spectrum

f_b : frequency break

- PSP measurement of magnetic power spectrum (Chen et al 2020)
 - Turbulence evolution down to 0.17 AU
 - Frequency break f_b which separates $1/f$ range and inertial range turbulence

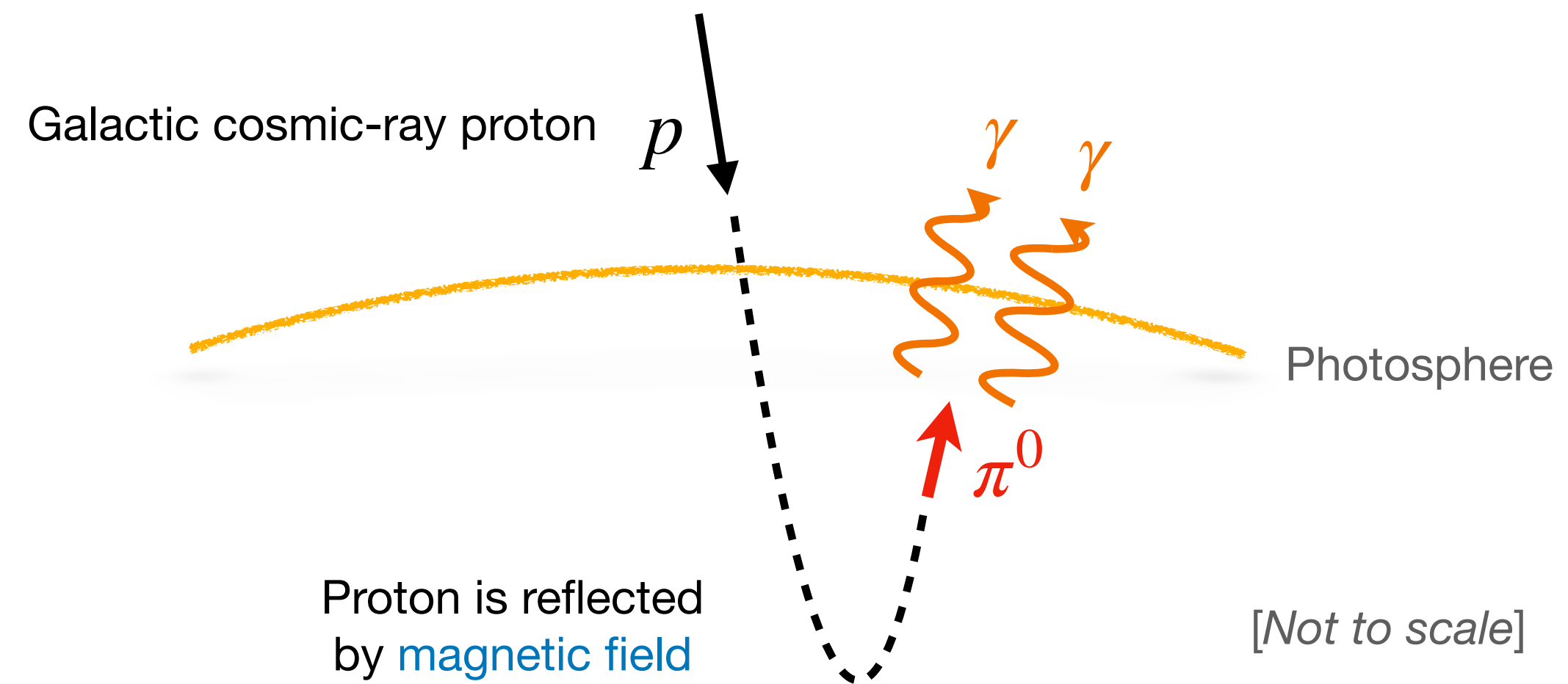


Cosmic-Ray Energy Spectrum at 0.1 AU

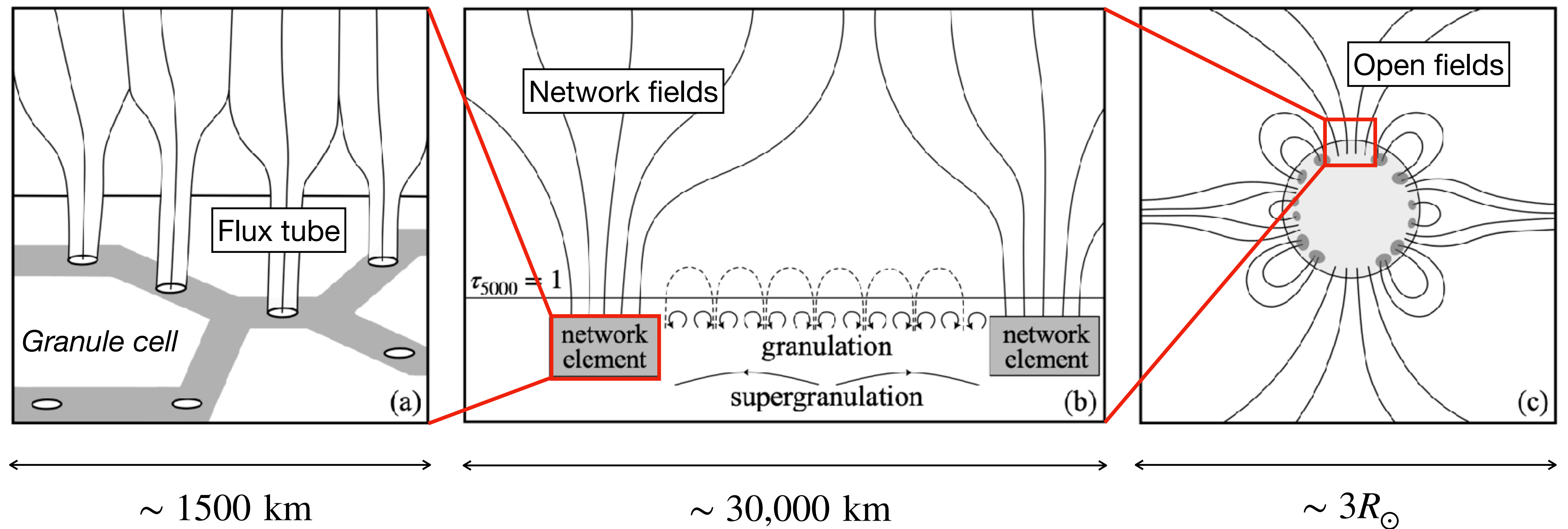


Modulation in the inner heliosphere is modest
 $\approx 10\%$ reduction of intensity from 1 AU to 0.1 AU

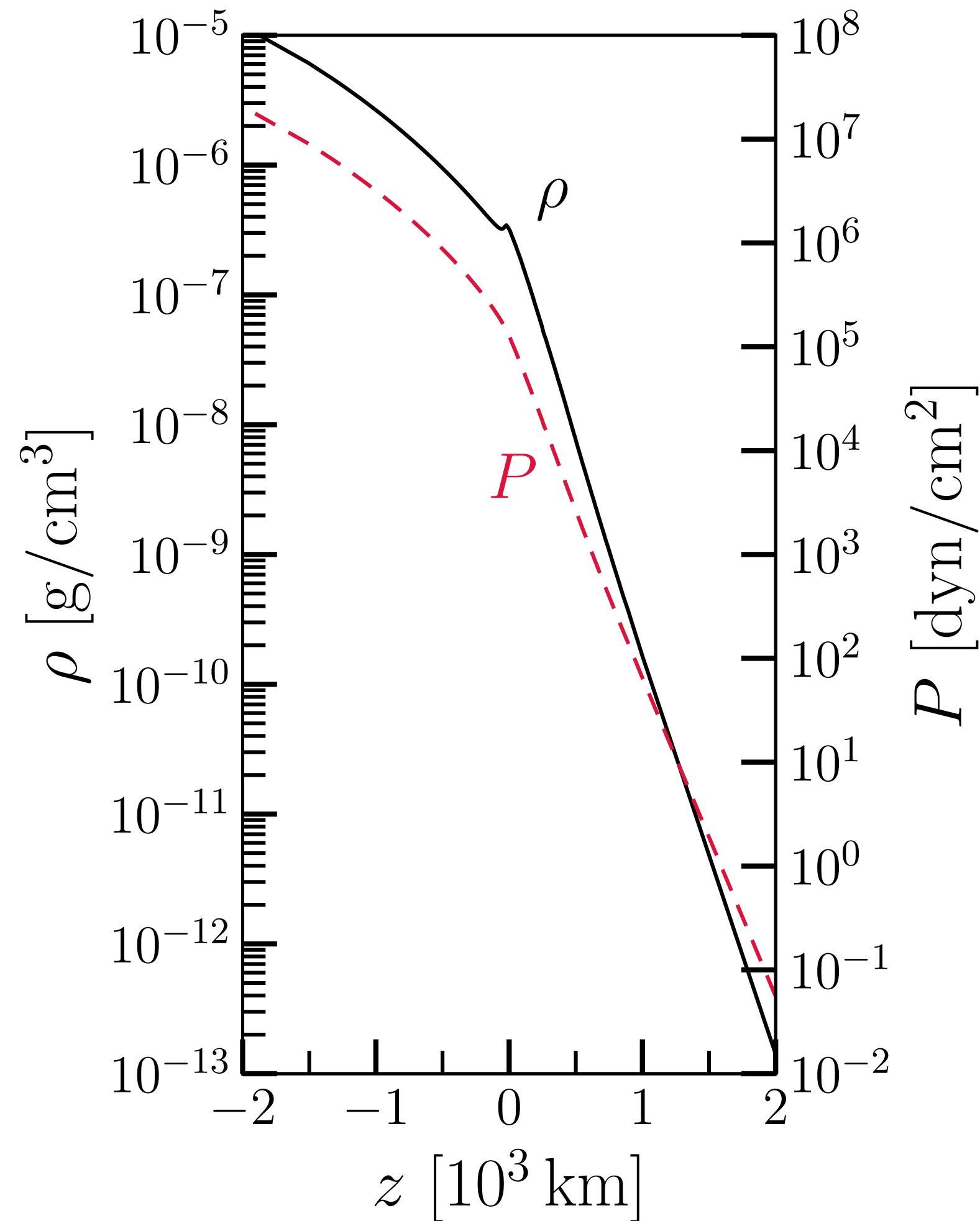
Stage 2: Particle reflection in flux tube & flux sheet



Overview of Coronal-hole Open Field Lines & Magnetic Network Fields (Quiet Photosphere Region)



Depths of Interest



- Magnetic field structure is multi-scale — need to identify the Depth of interest for gamma-ray production

- Estimates of proton GCR absorption in the Sun

- One absorption from pp interaction

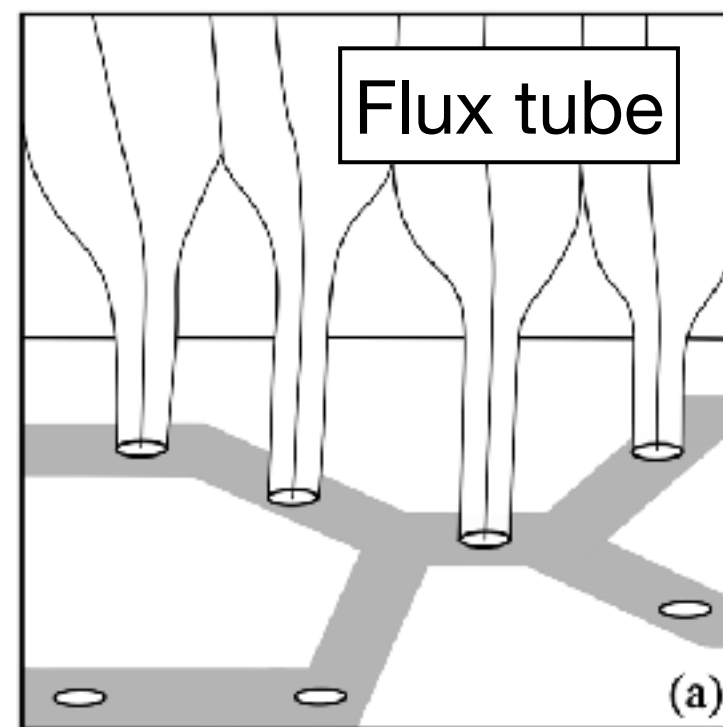
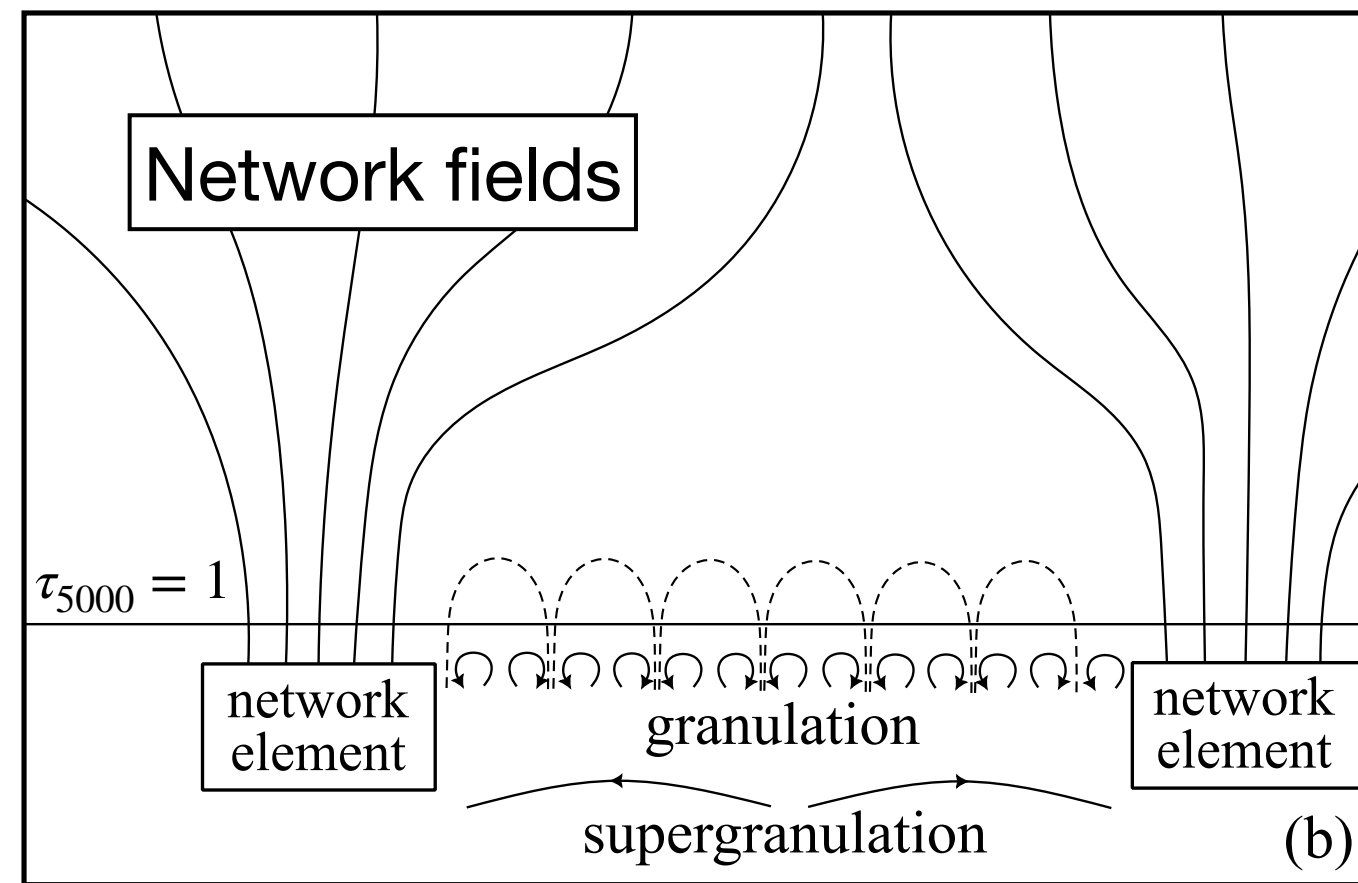
$$\int n_{\text{gas}}(z) \sigma_{pp} dz \sim 1$$

- pp interaction occurs within ~ few 100 km below solar surface.

- Surface ($z=0$) is defined as $\tau_{500\text{ nm}} = 1$

- Gamma rays are emitted in *photosphere* and *uppermost convection zone*

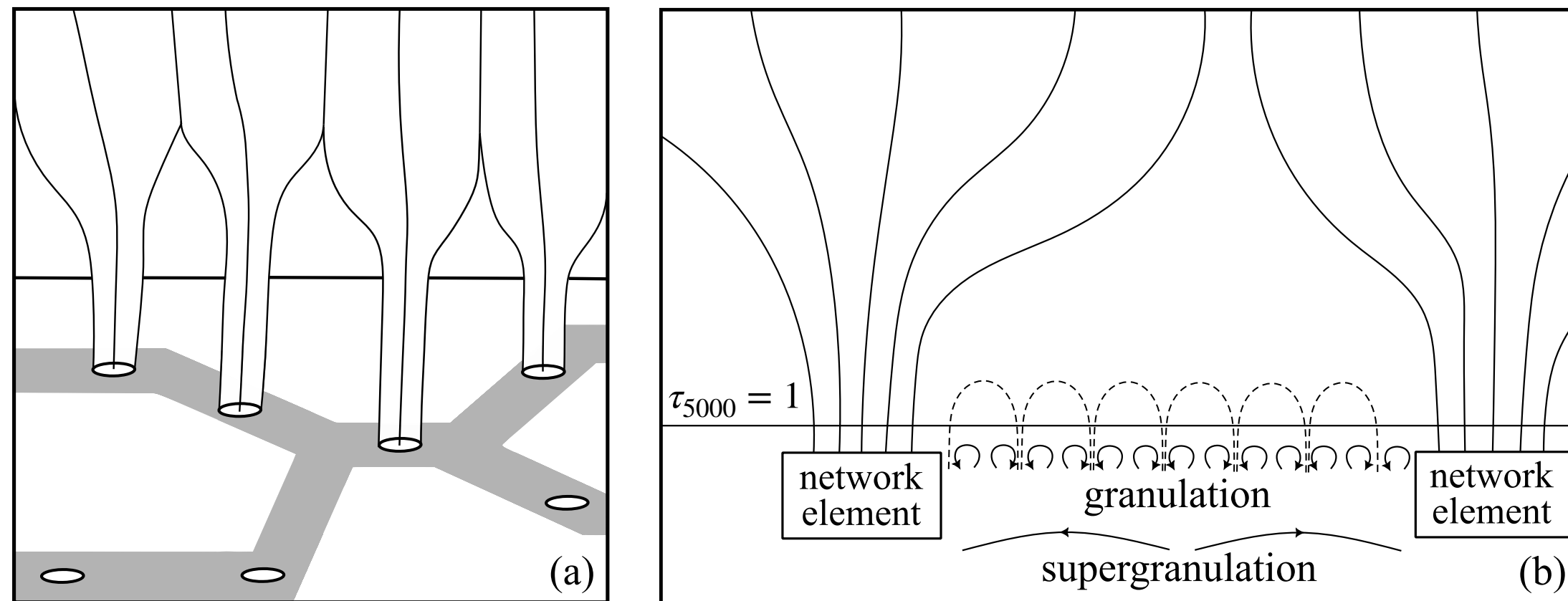
Our Model Assumptions



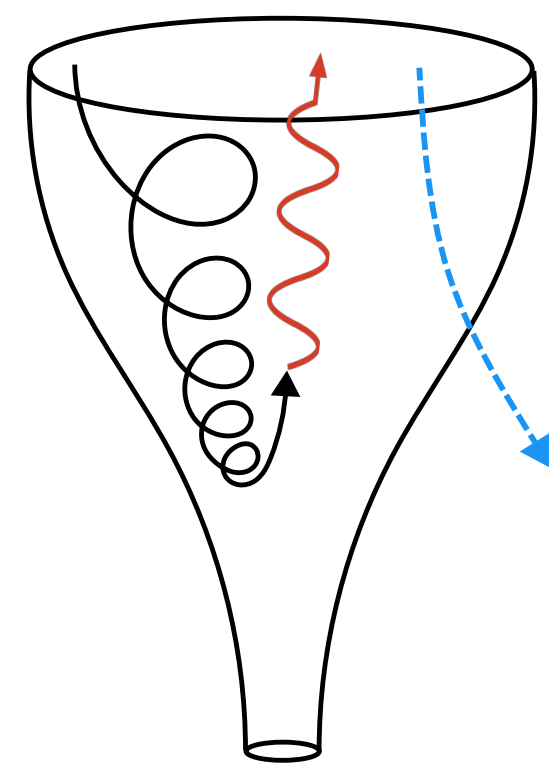
- GCRs propagate along open field lines, entering network elements
- Simulation starts at the merging height of tubes (at $z=1600$ km)
 - i.e., we consider chromosphere, photosphere, uppermost convection zone
- GCR intensity taken from AMS + CREAM measurement at 1 AU
 - Using Parker Solar Probe result on magnetic power spectrum, GCR flux reduction is $\lesssim 10\%$ from 1 AU to 0.1 AU (see **JTL** et al 2022: ApJ **937** 27)
 - Solar modulation from 1 AU to solar surface is not considered
- Inject GCRs into tube isotropically
 - Those high-energy GCRs passing through tube surface enters internetwork regions consisting of sheets

JTL, Beacom, Griffith, Peter 2023
(arXiv: 2307.08728)

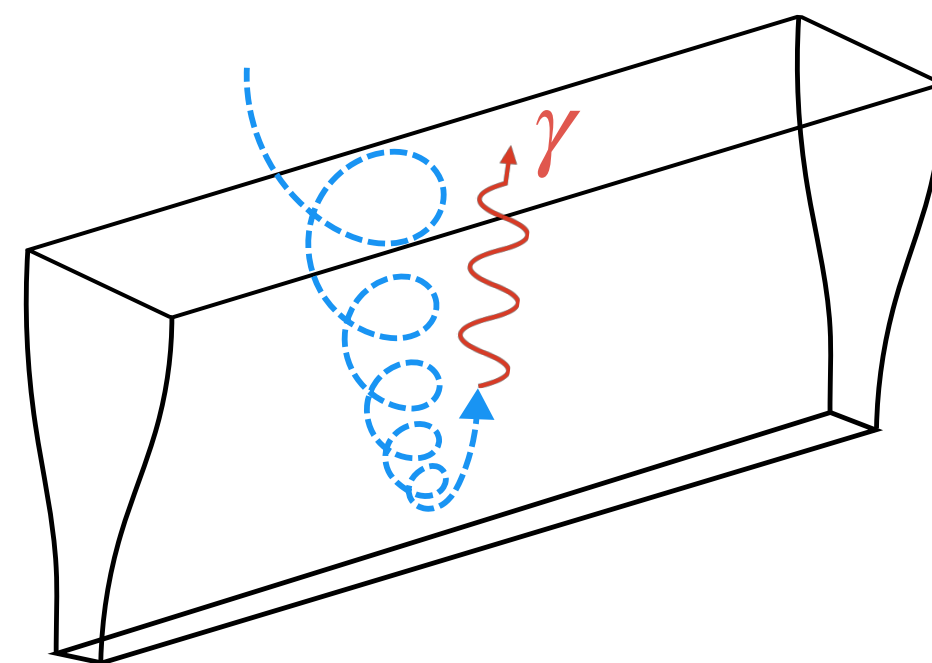
Schematics of Our Model: Flux Tube + Flux Sheet



- One flux tube and one flux sheet
 - Tube represents network element
 - Sheet represents granule sheet



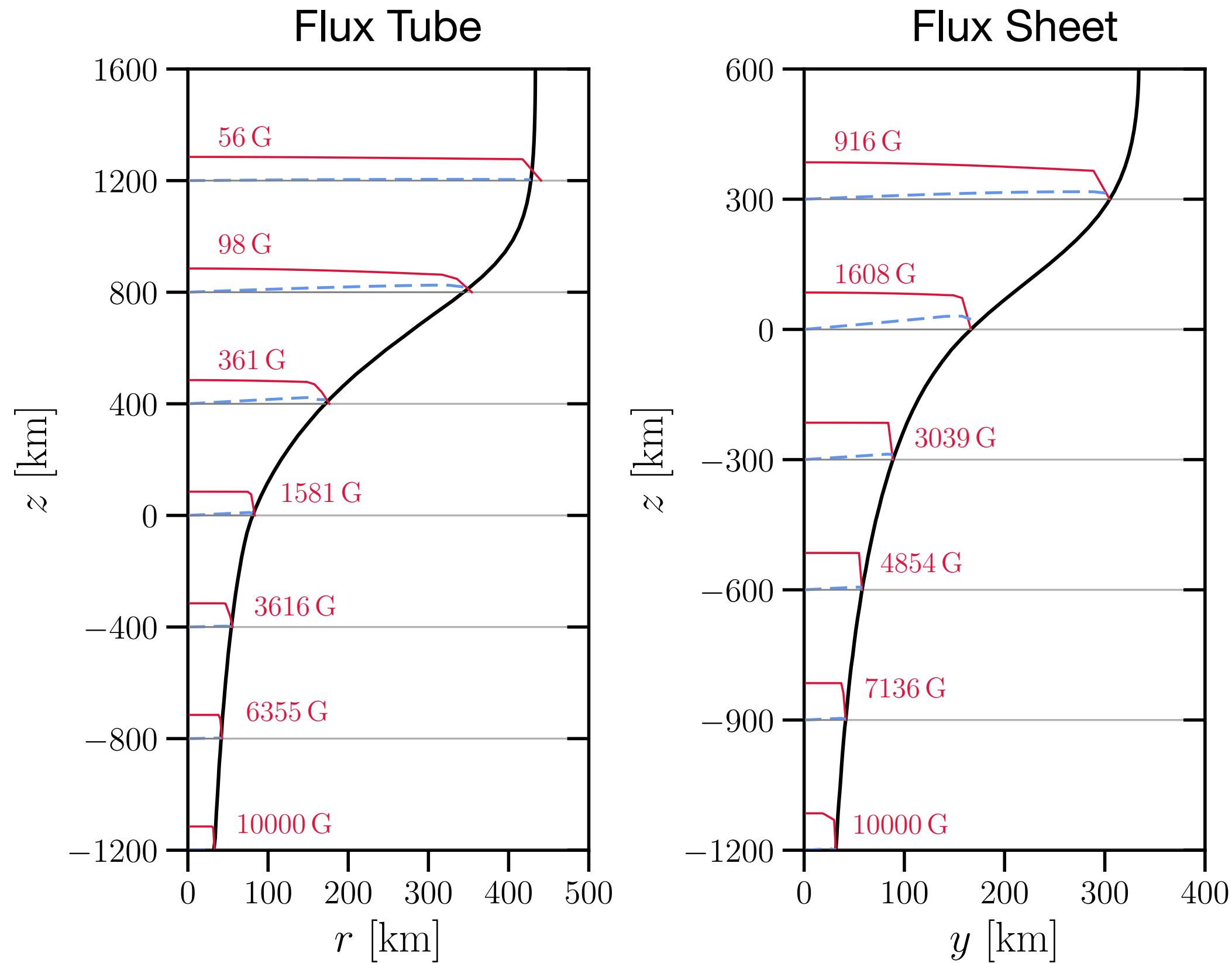
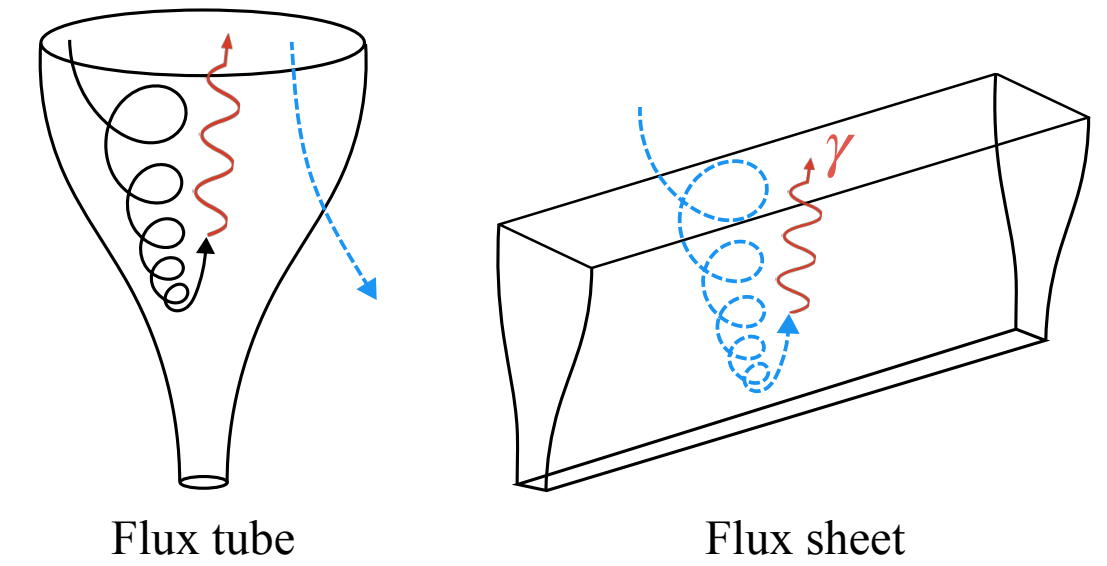
Flux tube



Flux sheet

- Particles are reflected via **magnetic bottle (magnetic mirroring) effect**
 - Increasing B makes pitch angle approaching 90°
 - Radial field imparts a kick at 90°
 - Particle starts spiraling upward

Magneto-Hydrostatic Equilibrium



- Magneto-hydrostatic equilibrium with the surrounding gas
- Following Grad-Shafranov equations

○ Flux tube:

$$\frac{\partial^2 \Psi}{\partial r^2} - \frac{1}{r} \frac{\partial \Psi}{\partial r} + \frac{\partial^2 \Psi}{\partial z^2} = -4\pi r J$$

$$B_r = -\frac{1}{r} \frac{\partial \Psi}{\partial z}, \quad B_z = \frac{1}{r} \frac{\partial \Psi}{\partial r}, \quad B_\phi = 0$$

○ Flux sheet:

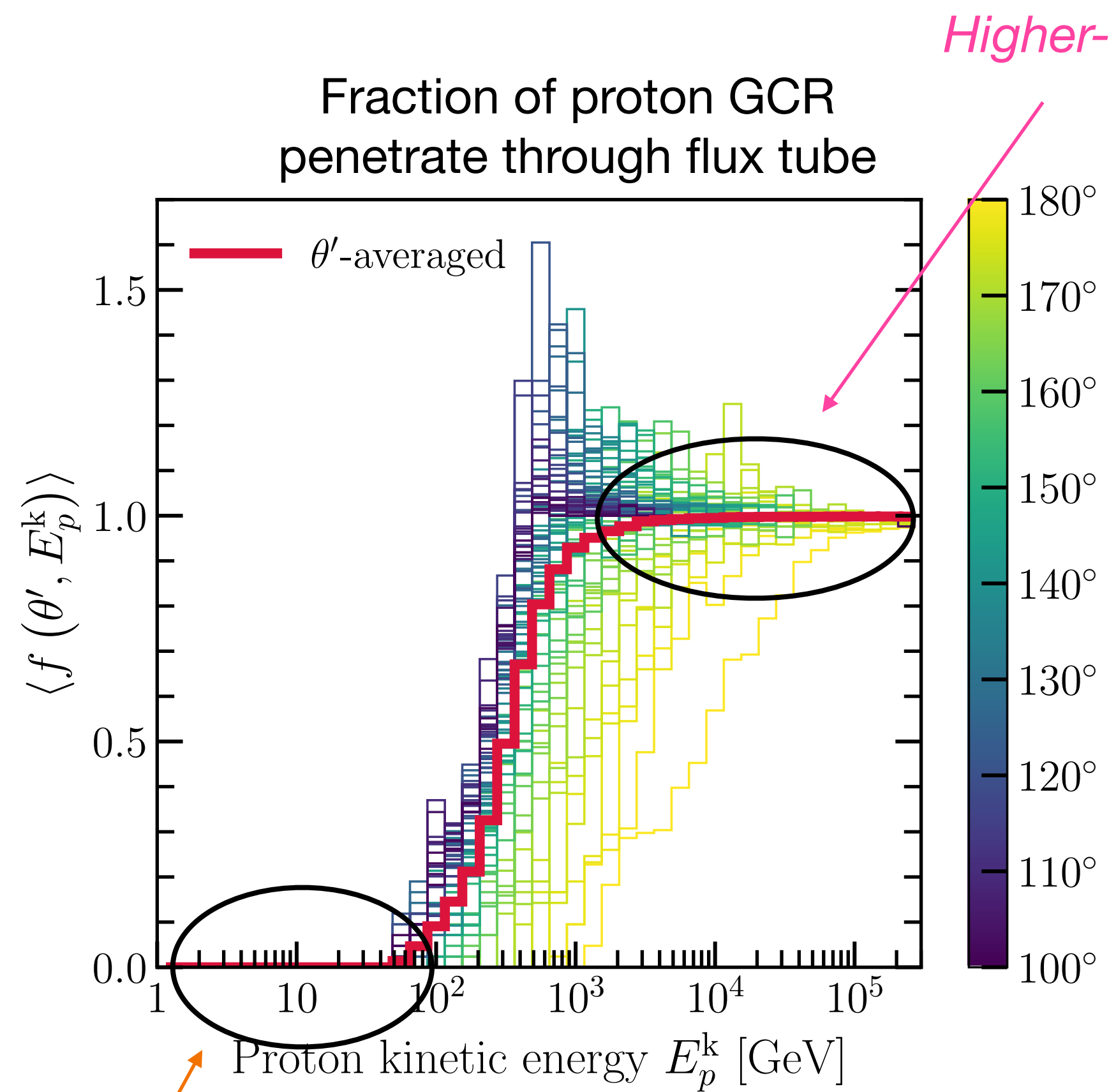
$$\frac{\partial^2 \Psi}{\partial y^2} + \frac{\partial^2 \Psi}{\partial z^2} = -4\pi J$$

$$B_y = -\frac{\partial \Psi}{\partial z}, \quad B_z = \frac{\partial \Psi}{\partial y}, \quad B_x = 0$$

- Internal magnetic flux structure is critical for magnetic bottle effect!

JTL, Beacom, Griffith, Peter 2023
(arXiv: 2307.08728)

Angular Distribution of Proton GCR Escaping Flux Tube

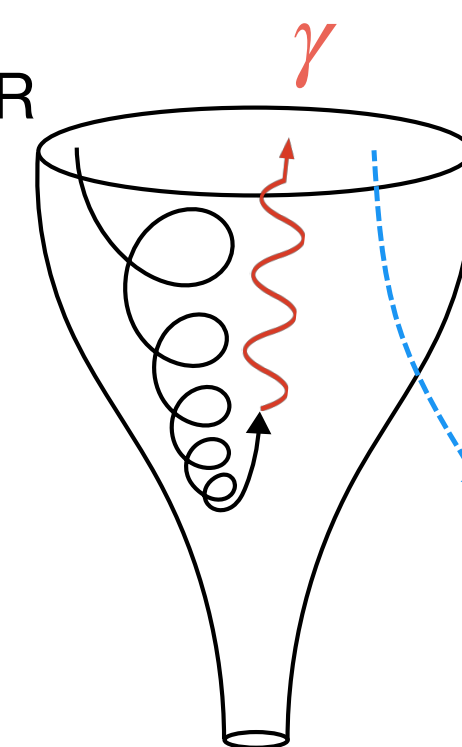


Higher-energy proton GCR passing through the flux tube, entering flux sheet

- Angular distribution: fraction of pGCR passing through flux tube, entering into flux sheet
- Low energy GCRs are tightly bounded by magnetic field lines in the tube
 - Cannot penetrating through tube
- High-energy GCRs are NOT bounded by magnetic field lines in the tube
 - Penetrating through tube, entering into internetwork regions (sheets)

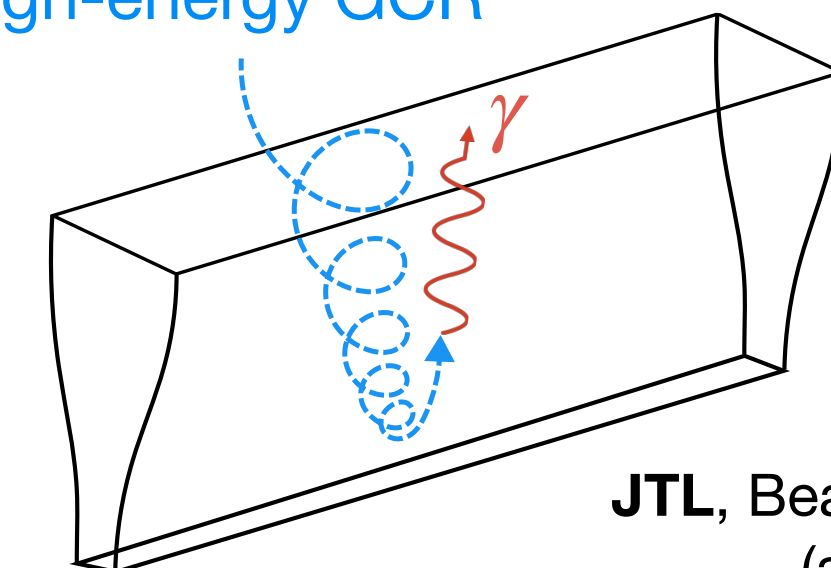
Lower-energy proton GCR bounded by the flux tube

Black:
Low-energy GCR



Flux tube

Blue:
High-energy GCR



Flux sheet

JTL, Beacom, Griffith, Peter 2023
(arXiv: 2307.08728)

Calculation of Gamma-Ray Emission

- Main gamma-ray production channel: $p + p \rightarrow p + p + \pi^0, \quad \pi^0 \rightarrow \gamma + \gamma$
- Gamma-ray flux

$$\frac{dN_\gamma}{dE_\gamma} = \int_{\Omega_0} \int_{E_\gamma}^{\infty} \int_0^{\bar{\chi}_p} F_\gamma(E_\gamma, E_p) \Phi_p(E_p) \cos \theta_0 \frac{dP_{\text{abs}}(\chi_p, E_p)}{d\chi_p} \zeta(\mathbf{r}) d\chi_p \frac{dE_p}{E_p} d\Omega_0$$

$$\frac{dN}{dE_\gamma} \sim (\# \text{ of } \gamma \text{ per interaction}) \times (\text{GCR flux}) \times (\text{GCR absorption prob.}) \times (\gamma \text{ transmission prob.})$$

Gamma-ray yield only available for $E_\gamma \gtrsim 3 \text{ GeV}$
In the literature

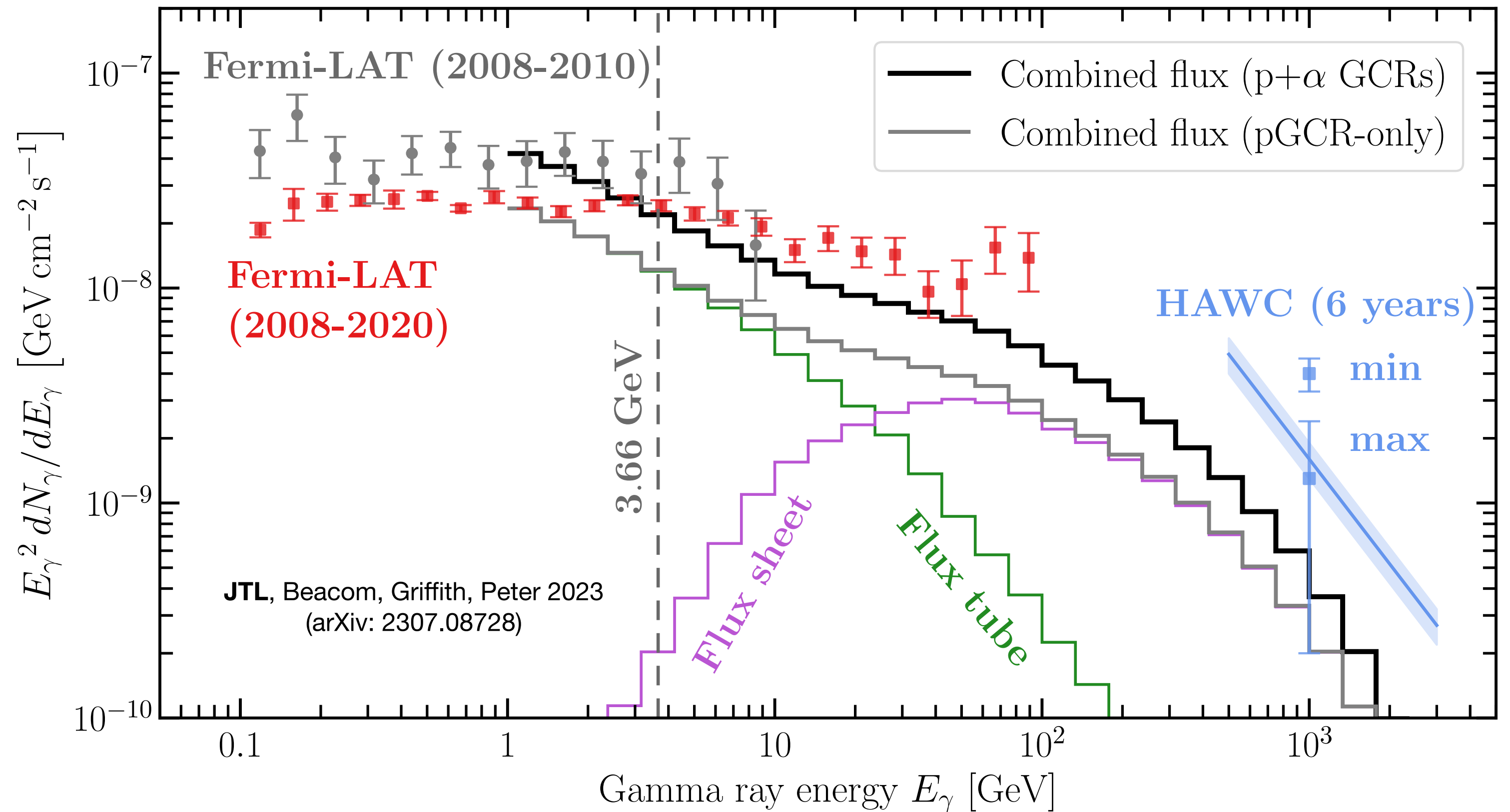
Kelner et al 2006
(PRD 74, 034018)

Numerical calculation of GCR trajectory

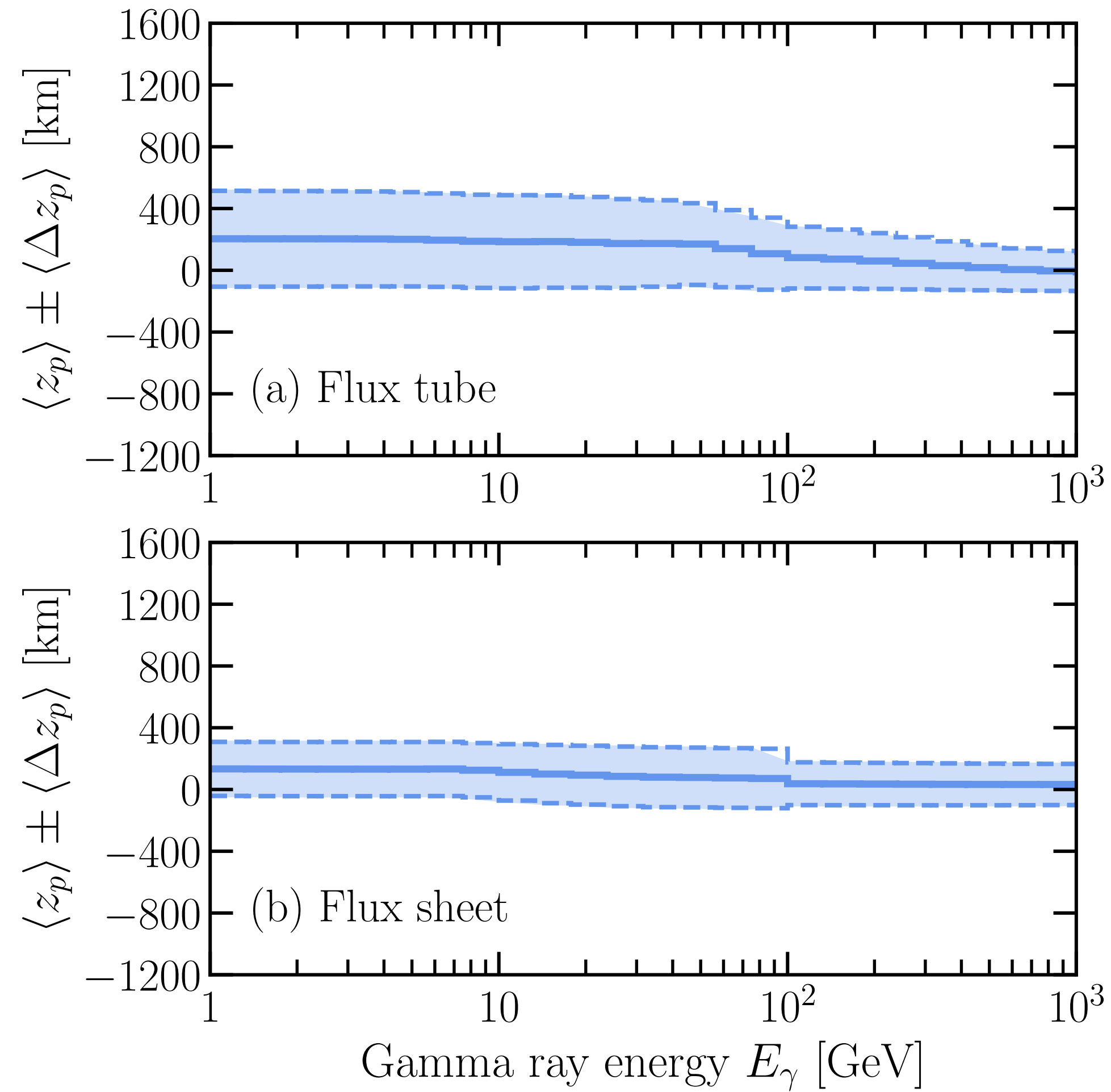
Gamma rays transmitted out from the solar gas

Our Result: Gamma-Ray Spectrum

- Lower-energy ($\lesssim 10$ GeV) gamma rays are produced from **flux tube (forming the network element)**
- Mid-energy ($1 \text{ GeV} \lesssim E_\gamma \lesssim 100 \text{ GeV}$) gamma rays are produced from the combination of **flux tube** and **flux sheet**
 - Convective cell plays critical role!
- Higher-energy ($\gtrsim 100 \text{ GeV}$) gamma rays produced from **flux sheet (between granular convective cells)**
 - GCR isotropically bombard internetwork regions



Average Emission Height

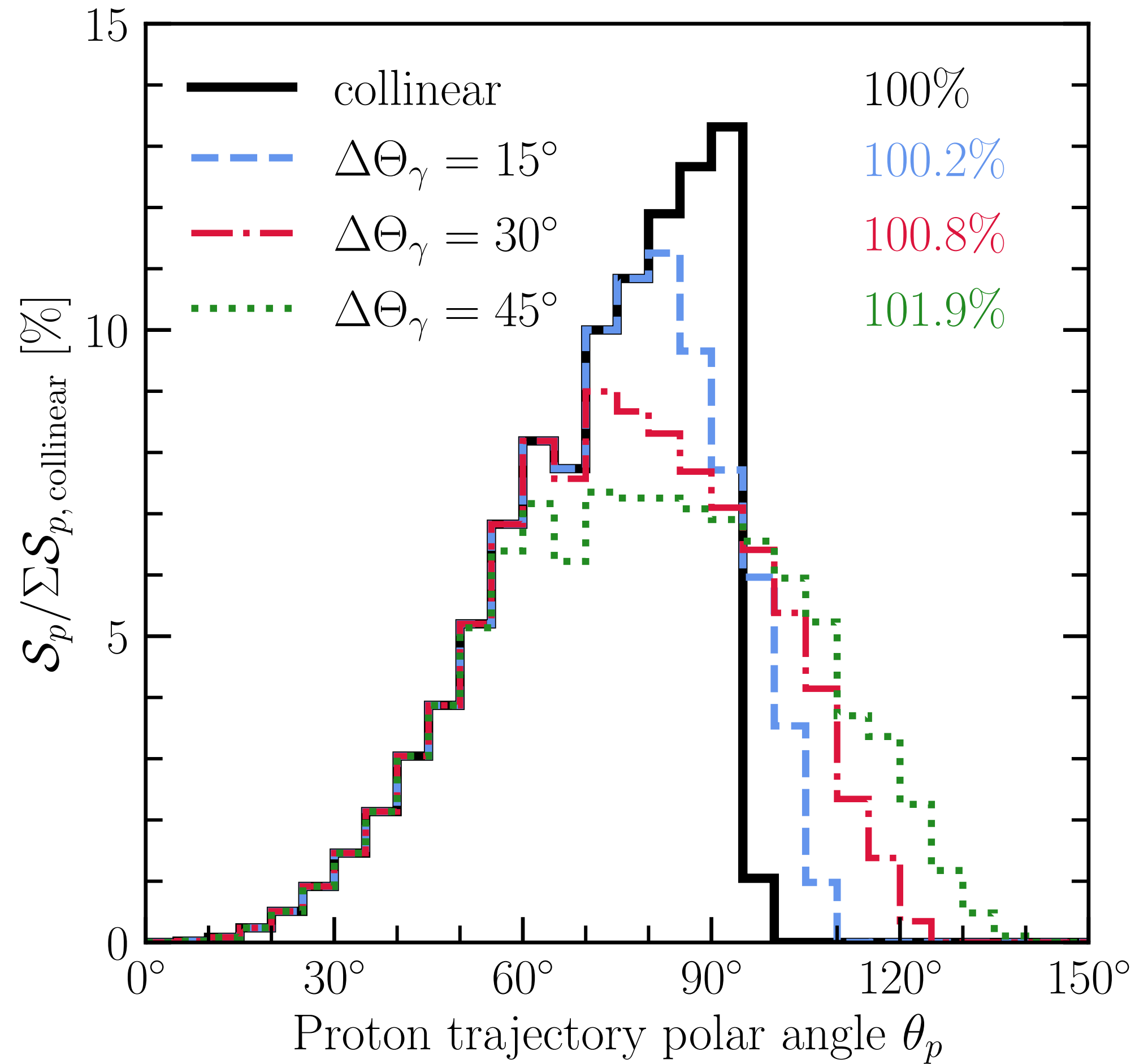


- Emission occurs primarily in the chromosphere, photosphere, and upper-most convection zone (~100 km)

Conclusions and Outlook

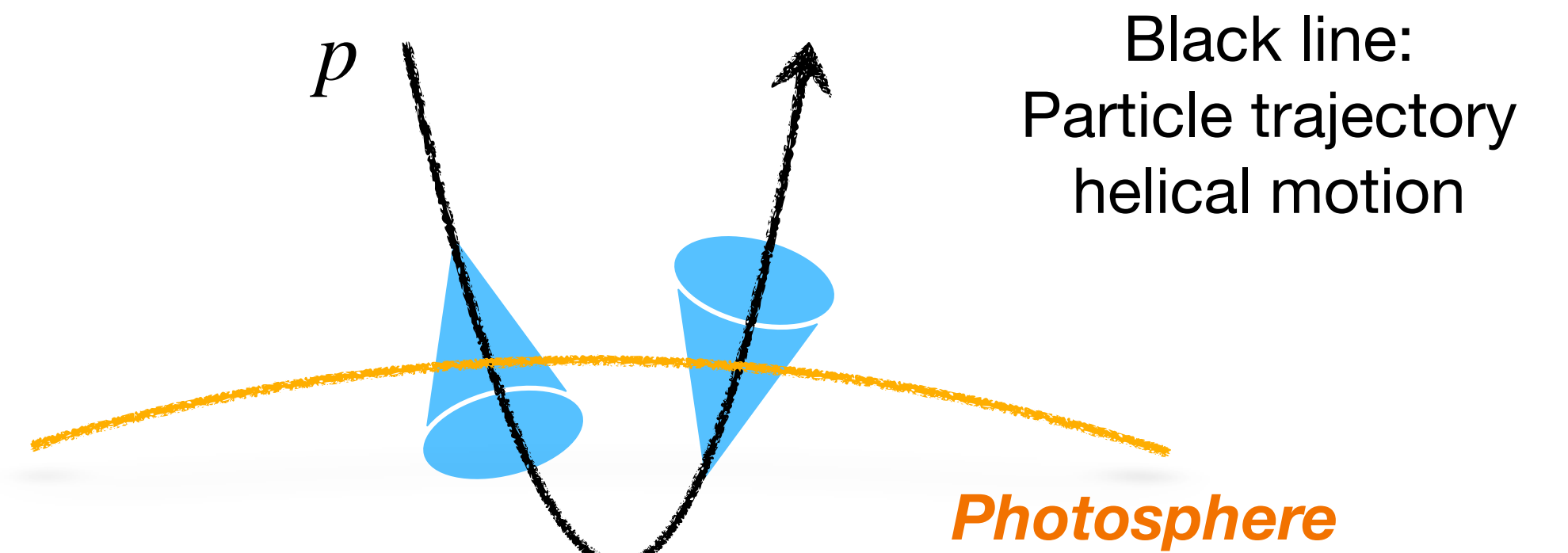
- A simple model consisting of **one tube** and **one sheet**
 - Gamma-ray observation data is explained reasonably well (within a factor 2)
 - Ineffectiveness of capturing high-energy GCRs causes the steep gamma spectrum at \sim TeV (HAWC)
- What causes the **anti-correlation** between gamma-ray flux and solar cycle?
 - Coronal holes? Active regions? Small-scale dynamo? GCR transport?
- How does **turbulence** from the **convective flow** affect GCR transport in the photosphere and uppermost convection zone?

Finite-Sized Emission Cone (for each pp interaction)

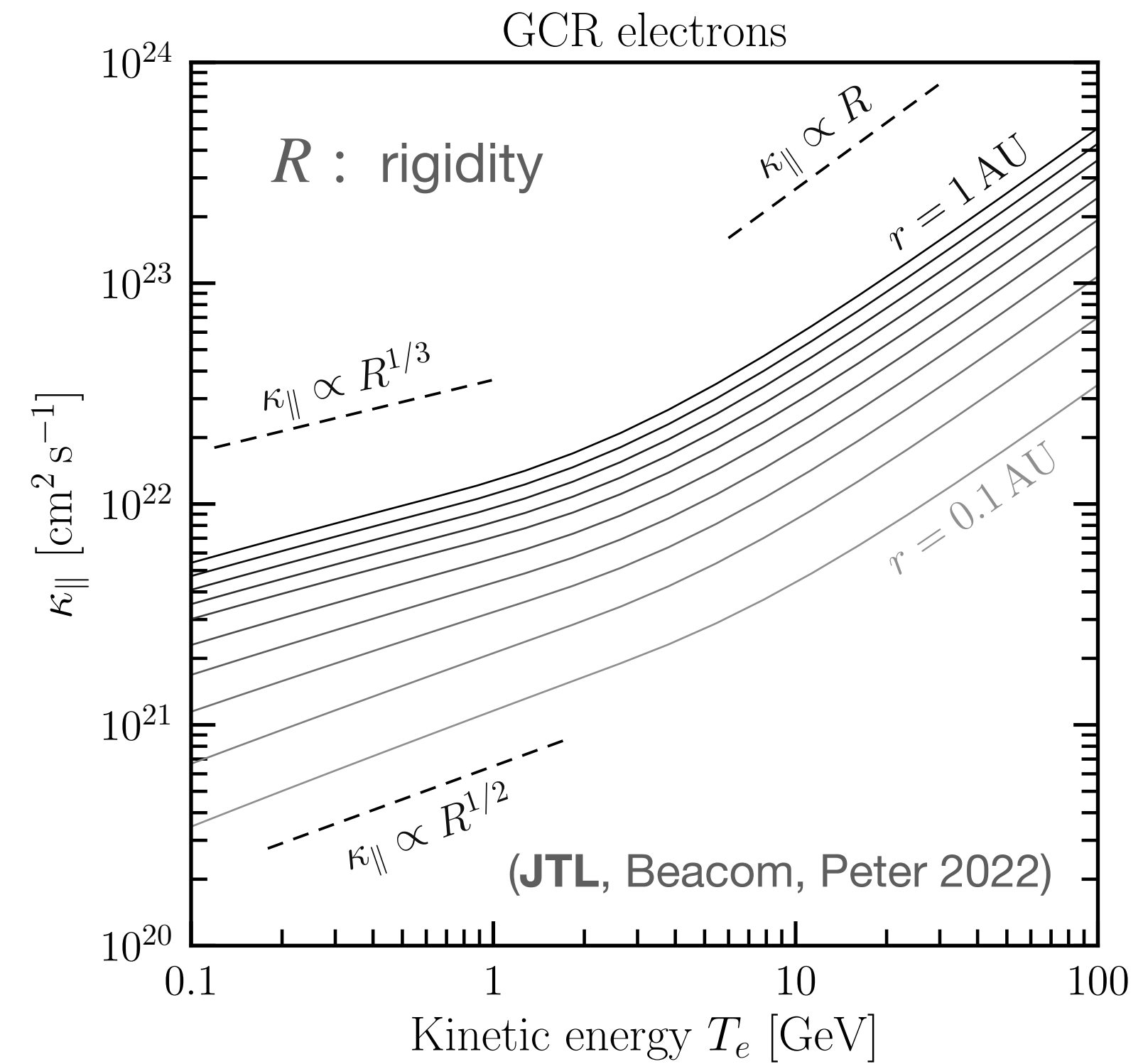
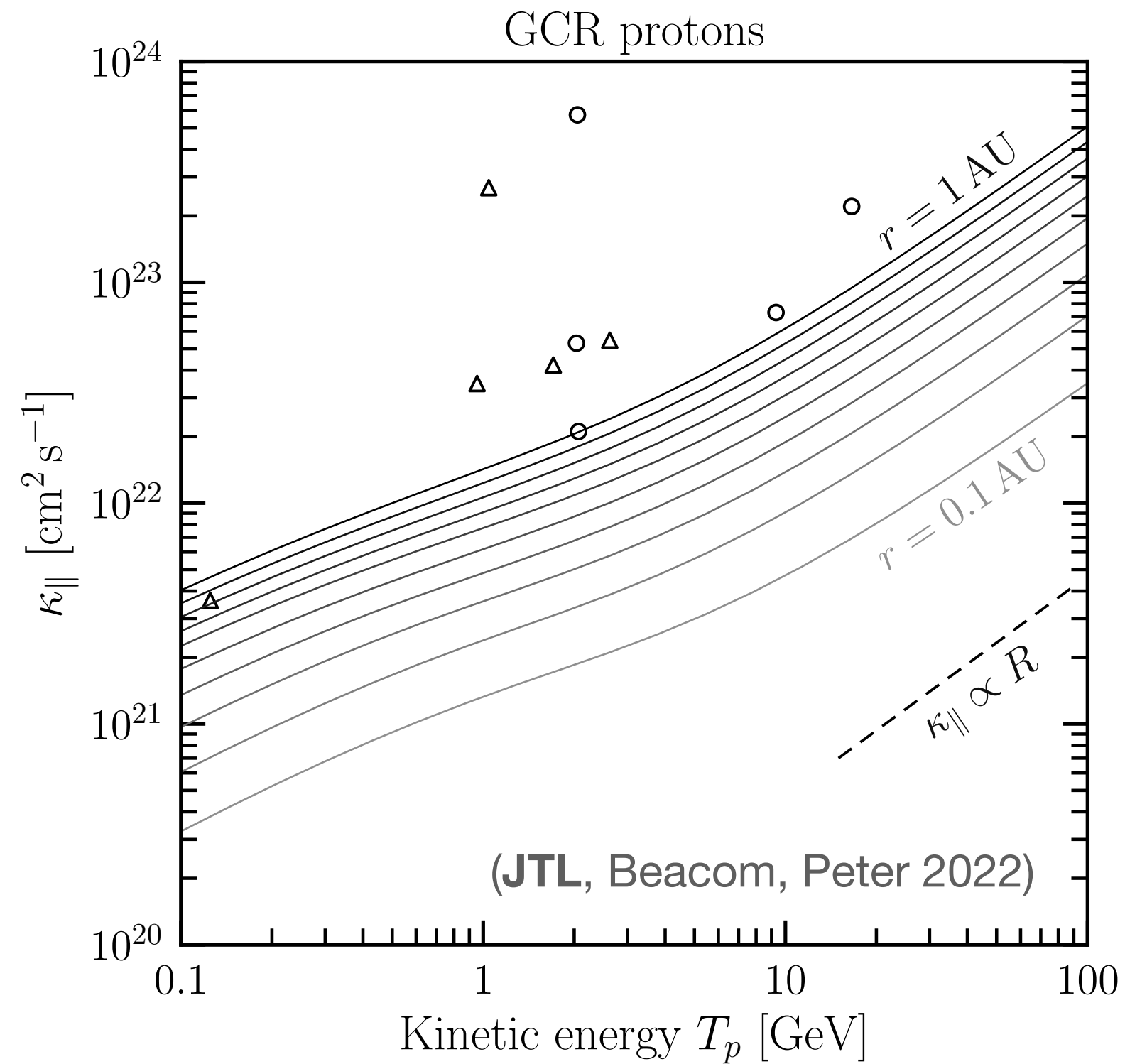


$$\mathcal{S}_p = \int_0^{\bar{\chi}_p} \frac{dP_{\text{abs}}(\chi_p, E_p)}{d\chi_p} \zeta(\mathbf{r}) d\chi_p$$

= proton GCR absorption probability
 × gamma absorption probability



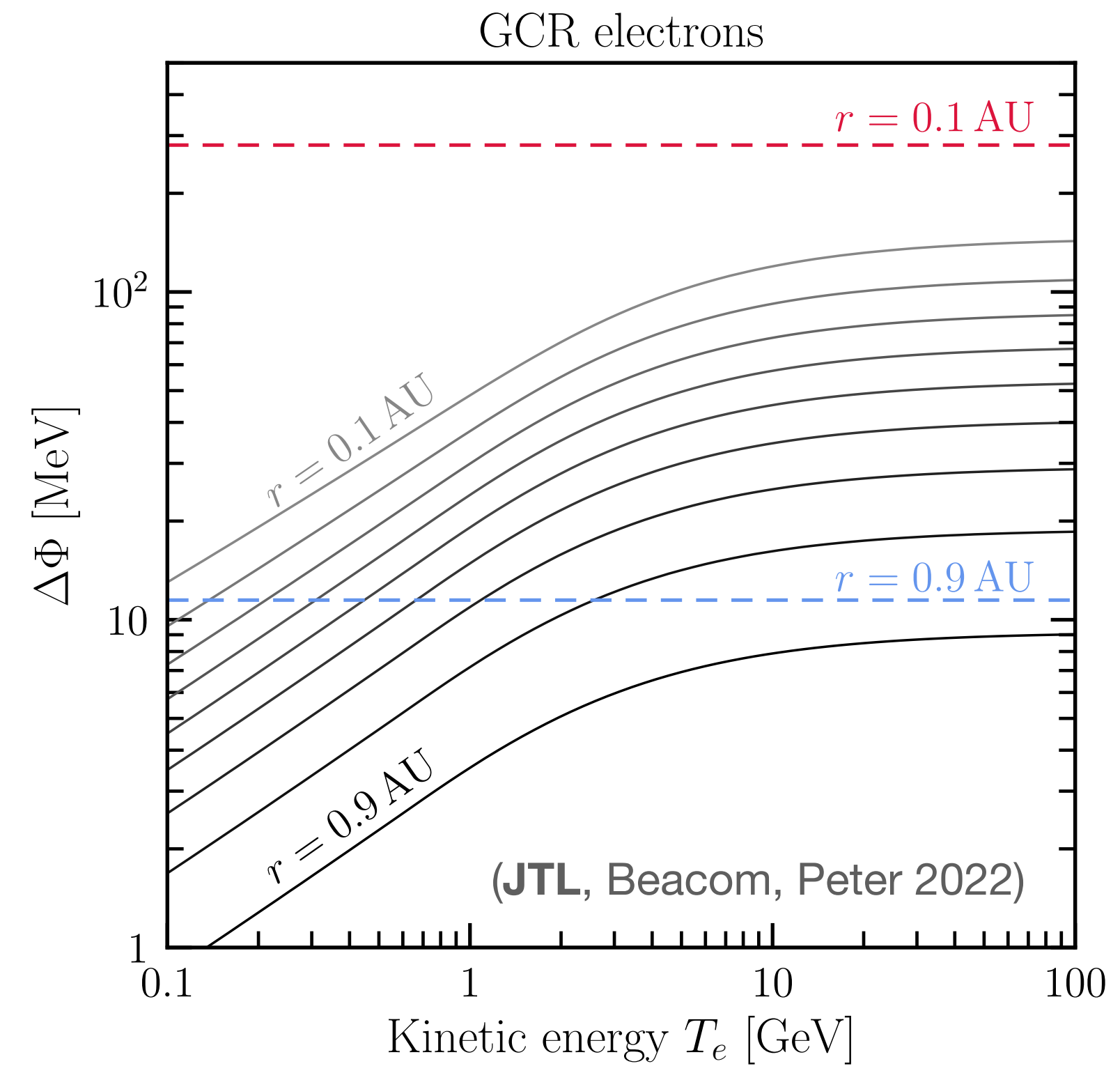
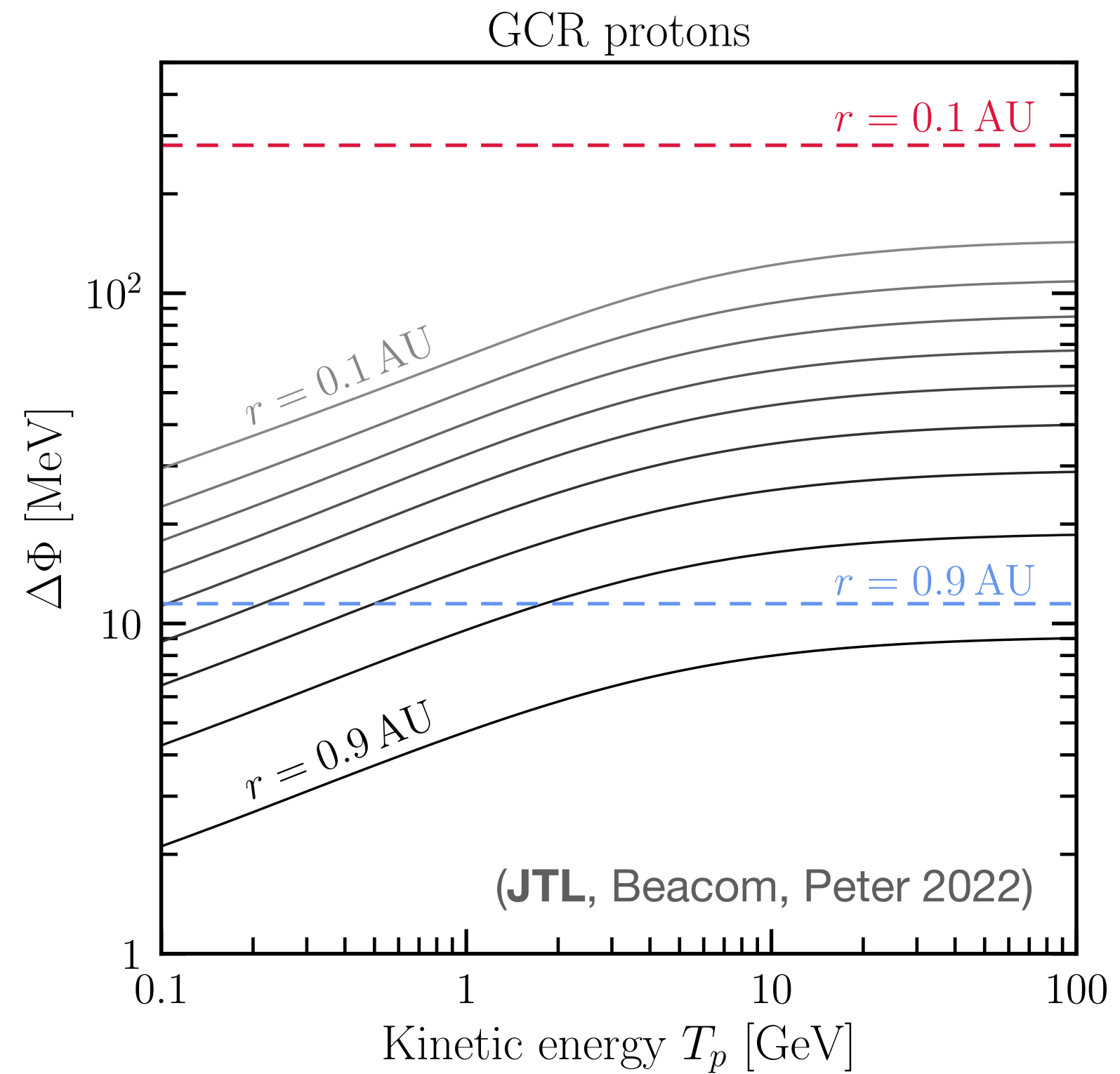
Diffusion Coefficients



Circle and triangle: measurements of CR proton, from Palmer 1982

Measured mean free path is approximately 2 times higher than QLT result, known as Palmer consensus

Modulation Potential Energy



Small modulation potential increase for $E_{\text{kin}} \lesssim 10$ GeV

Magnetic spectrum ($1/f$ v.s. inertial range) matters